Mass in Particle Physics QCD and the Mass of Hadrons

Thomas Mannel

Theoretische Physik I, Universität Siegen



Schule für Astroteilchenphysik 2014

A B A B A
 A
 B
 A
 A
 B
 A
 A
 B
 A
 A
 B
 A
 A
 B
 A
 A
 B
 A
 A
 B
 A
 A
 B
 A
 A
 B
 A
 A
 B
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A

Contents



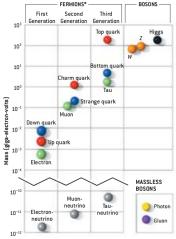
2 Dimensional Analysis for Pedestrians

Mass from dimensional Transmutation

くロト (過) (目) (日)

Introduction / Motivation

Masses of the fundamental fermions



T. Mannel, Siegen University

Mass in Particle Physics: Lecture 2

(日)((日))

ъ

- The proton consists of up and down quarks
- We have [PDG]

 $m_u = 2.3^{+0.7}_{-0.5} \,\mathrm{MeV}$ and $m_d = 4.8^{+0.5}_{-0.3} \,\mathrm{MeV}$

How can the proton mass be 938 MeV?

- Almost the full mass must be "binding energy"
- Certainly not a loosely-bound system
- The mass of the quarks seems to be negligible, so ...

How can the proton mass for vanishing quark masses?

ヘロン ヘアン ヘビン ヘビン

Quantum Field Theory in a Nutshell

We have considered classical field theory for the SM
QCD is part of the SM:

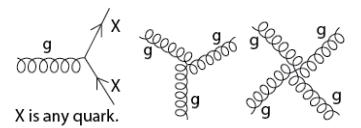
$$\mathcal{L}_{ ext{QCD}} = \sum_{m{q}=u, m{d}, m{s}} ar{m{q}} i m{D} m{q} - rac{1}{4} m{F}^a_{\mu
u} m{F}^{\mu
u, m{a}}$$

where we look only at the light quarks q=u,d,s and the gluons ${\cal F}^a_{\mu
u}$

- There is no dimensional quantity in this Lagrangian
- ... so how can a mass emerge?

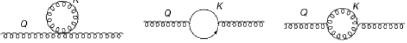
ヘロト ヘ戸ト ヘヨト ヘヨト

- (Perturbative) Quantum Field Theory
- Feynman Rules



- Feynman Digramms: Define your external lines and draw all possible diagrams with a given number of vertices
- This translates into a mathematical expression for the quantum-mechanical amplitude.





- The "loop-momentum" can be infinitely high
- Integrand does not cut-off the high momentum modes
- These contributions need "renormalization"
- Renormalization is a general feature of QFT

The Essence of Renormalization

- Physical parameters (masses and couplings) are obtained form subtracting a "counter term contribution" from the "bare parameter"
- Renormalization requires to fix the parameters of the theory by experimental input.
- Example: The QCD coupling is determined by performing a scattering experiment at a high scale μ

$$\sigma(\mu) \sim \frac{\alpha_s(\mu)}{\mu^2}$$

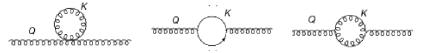
- This defines a "running coupling"
- Perturbative expansion in $\alpha_s(\mu)$
- This always introduces a "renormalization scale" µ

Running Coupling

• The running is defined by the β function

$$\mu \frac{\boldsymbol{d}}{\boldsymbol{d}\mu} \alpha_{\boldsymbol{s}}(\mu) = \beta(\alpha_{\boldsymbol{s}}(\mu))$$

 the β function is (perturbatively) calculated from the diagrams



Result known up to four loops

< ∃ > <

One loop result

$$\beta(\alpha_s) = \left(\frac{2n_f}{3} - 11\right)\frac{\alpha_s^2}{2\pi} + \mathcal{O}(\alpha_s^3) = \beta_0 \alpha_s^2 + \mathcal{O}(\alpha_s^3)$$

General solution for the running coupling

$$\frac{d\mu}{\mu} = \frac{d\alpha_s}{\beta(\alpha_s)} \qquad \ln\left(\frac{\mu}{\mu_0}\right) = \int_{\alpha_s(\mu_0)}^{\alpha_s(\mu)} \frac{d\alpha}{\beta(\alpha)}$$

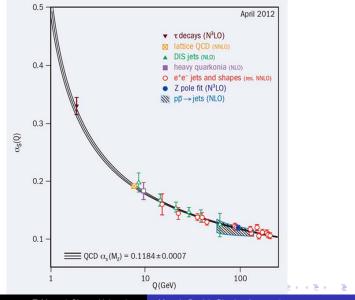
One loop running

$$\beta_0 \ln\left(\frac{\mu}{\mu_0}\right) = \frac{1}{\alpha_s(\mu_0)} - \frac{1}{\alpha_s(\mu)} \qquad \alpha_s(\mu) = \frac{\alpha_s(\mu_0)}{1 - \beta_0 \ln\left(\frac{\mu}{\mu_0}\right)}$$

・ロン ・雪 と ・ ヨ と

Quantum Field Theory in a Nutshell

Dimensional Analysis for Pedestrians Mass from dimensional Transmutation



T. Mannel, Siegen University

Mass in Particle Physics: Lecture 2

Dimensional Analysis for Pedestrians

Consider a dimensionless quantity R (such as $s\sigma(s)$)

and pretend for a moment there were no renormalization !

Perturbative expansion in QCD

$$\boldsymbol{R} = \boldsymbol{R}(\alpha_{\boldsymbol{s}}) = \sum_{n=0}^{\infty} \boldsymbol{r}_n \alpha_{\boldsymbol{s}}^n$$

 In massless QCD (no scale!) this cannot depend on any momentum transfer q²

$$R(q^2)
ightarrow R\left(rac{q^2}{\Lambda^2}
ight), \qquad r_n ext{ and } lpha_s ext{ dimensionless}$$

but there is no Λ in (finite) QCD.

→ E > < E >

- If everything would be finite, naive dimensional analysis would hold
- Conclusion: at sufficiently high scales $\sigma(s) \sim 1/s$
- However: α_s is running, i.e. it depends on a scale
- ", but which is arbitrary

How can we get something from nothing?

ヘロト ヘアト ヘビト ヘビト

Mass from dimensional Transmutation

- The renormalization scale is still arbitrary Nothing can depend on it
- How to get a "renormalization invariant" mass scale?
- Construct $\Lambda_{QCD}(\mu, \alpha_s(\mu))$

$$\Lambda_{\text{QCD}}(\mu, \alpha_{s}(\mu)) = \mu \exp\left(-\int_{\alpha_{s}(\mu_{0})}^{\alpha_{s}(\mu)} \frac{d\alpha}{\beta(\alpha)}\right)$$

Show that

$$\mu \frac{d}{d\mu} \Lambda_{\text{QCD}}(\mu, \alpha_{s}(\mu)) = \mathbf{0} = \left(\mu \frac{\partial}{\partial \mu} + \beta(\alpha_{s}) \frac{\partial}{\partial \alpha_{s}}\right) \Lambda_{\text{QCD}}(\mu, \alpha_{s})$$

ヘロン 不通 とくほ とくほ とう

One loop result

$$\Lambda_{\text{QCD}}(\mu, \alpha_s) = \operatorname{const} \mu \exp\left(-\frac{1}{\beta_0 \alpha_s}\right)$$

(the constant is usually put to unity)

- Value of Λ_{QCD} depends on the scheme in which renormalziation is performed
- Usual scheme is "modified minimal subtraction" MS

$$\Lambda_{\rm QCD,\,\overline{MS}} = (340\pm8)~{
m MeV}$$

(for three quark flavors)

QCD generates a mass scale of a few hundred MeV!

<ロト < 同ト < 回ト < 回ト = 三

A few remarks:

- Renormalizable theories always generates their mass scale, even if you start out massless!
- this looks like something from nothing, but there is the input of running α_s
- This scale is nonperturbative: There is no Taylor series of exp(1/x)
- The dimensional analysis has to be modified due to the presence of this scale!

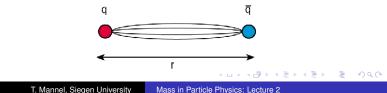
ヘロン 人間 とくほ とくほ とう

- Trivial example: Coulomb Potential
- Maxwell Theory is scaleless, no renormalization
- Potential as a function of r between static charges: Dimensional analysis: Energy ~ inverse length, thus

$$V(r) = rac{\mathrm{const}}{r}$$

- Pure-Glue QCD: scaleless, but needs renorm.:
- Potential between two static color charges:

$$V(r) = \frac{1}{r} f\left(\frac{r}{\Lambda_{\rm QCD}}\right) \to \sigma r$$



Intermediate Summary

- Massless QCD generates a mass scale due to renormalization
- Even in massless QCD we thus expect to have a non-vanishing proton mass!
- Explains the mismatch between the light quark masses (emerging from the Higgs mechanism) and the large proton mass,

くロト (過) (目) (日)

More on Masses in QCD: Chiral Symmetry

• Massless QCD (for the three light quarks) has an $U(3)_L \times U(3)_R$ Symmetry

$$\mathcal{L} = \sum_{k} ar{q}_{k,L} i oldsymbol{D} q_{k,L} + \sum_{k} ar{q}_{k,R} i oldsymbol{D} q_{k,R} + ext{gluons}$$

- However, that symmetry is not seen in nature, only the diagonal symmetry $SU(2)_{L+R}$ is actually seen.
- ", corresponding to Isospoin or Flavour SU(3)
- If the symmetry were exact: no proton mass
- ASSUMPTION: SU(3)_L × SU(3)_R → SU(2)_{L+R} is a spontaneous symmetry breaking (SSB)

ヘロト ヘアト ヘビト ヘビト

- This SSB generates Goldstone particles: Pions and Kaons
- The "order parameter" is the quark condensate $\langle \bar{q}q \rangle = \langle 0 | \bar{q}q | 0 \rangle$ (corresponding to the Higgs VEV)
- ... this couples left an right handed components
- Pions and Kaons get (small) masses from explicit symmetry breaking from the quark mass terms:

$$\mathcal{L}_{\mathrm{mass}} = ar{q}_{L} \mathcal{M} q_{R} + \mathrm{h.c.} \qquad \mathcal{M} = \left(egin{array}{ccc} m_{u} & 0 & 0 \ 0 & m_{d} & 0 \ 0 & 0 & m_{s} \end{array}
ight)$$

ヘロト ヘアト ヘビト ヘビト

Relation for the pion and kaon mass

$$m_{\pi}^2 = -rac{\langle \bar{q}q \rangle}{3f_{\pi}^2}(m_u+m_d)$$
 $m_K^2 = -rac{\langle \bar{q}q \rangle}{3f_{\pi}^2}(m+m_s),$

with
$$m = \frac{1}{2}(m_u + m_d)$$

• From this we get

 $\langle \bar{q}q \rangle \sim (-0.23 \text{ GeV})^3$

イロト 不得 とくほ とくほ とう

E DQC

Back to massless QCD:

 Also the quark condensate has to originate from dimensional transmutation

$$\langle \bar{q}q \rangle = \text{const.} \times \Lambda_{\text{QCD}}^3$$

with a constant of order unity

• This is a consistent picture, however, difficult to quantify

ヘロン 人間 とくほ とくほ とう

1

Finally: The mass of the proton

- Start from the proton state at rest: $| {m
 ho}
 angle$
- This should be an eigenstate of the hamiltonian:

$$|H|
ho
angle = \int d^3ec x \ T^{00}(x) \ket{
ho} = M_{
m Proton} \ket{
ho}$$

with the energy-momentum tensor $T^{\mu
u}$

• From this we get

$$raket{m p}{\mathcal T}^{00}(0) raket{m p} = M_{
m Proton}$$

• In massless QCD *T*⁰⁰ is scaleless, so the mass can only come form dimensional transmutation

ヘロン ヘアン ヘビン ヘビン

- The proton mass does not vanish in the chiral limit, thus it is not proportional to the quark masses
- approximate relation (loffe)

$$M_{
m Proton}^3\sim 2(2\pi)^3 \langle ar q q
angle$$

which actually yields $M_{\rm Proton} \sim 1$ GeV.

イロト 不得 とくほ とくほ とう

э.

Summary on Lecture 2

- The proton mass (and all other hadron masses, except for the loins an kaons) has NOTHING to do with the Higgs
- Even in the massless limit, QCD is believed to have a non-trivial mass spectrum
- The mechanism is dimensional transmutation, which generates a scale
- The hadron mass spectrum is related to the breaking of the chiral symmetry of QCD

ヘロア 人間 アメヨア 人口 ア