

Fast Single-Top Cross Section Predictions for Hadron Colliders with the HATHOR Program

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15. Oktober 2014

Introduction

- The Top Quark

- Calculating Hadronic Cross Sections

Extraction of Partonic Cross Sections

- Method 1: Modification of PDFs

- Method 2: Exact Extraction

- Implementation in HATHOR

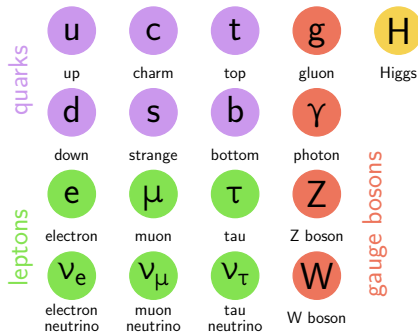
- Consistency Check

Practical Studies on the Production of Single Top-Quarks

Summary

The Top Quark in the Standard Model

- ▶ Standard model: fundamental interactions of elementary particles
- ▶ $m_t = 173.5 \text{ GeV}$, charge $+\frac{2}{3}e$
- ▶ Short lifetime: $5 \times 10^{-25} \text{ s}$
- ▶ Typical decay: $t \rightarrow Wb$
- ▶ Direct measurement of $|V_{tb}|^2$
- ▶ **Suitable laboratory** for new physics: e.g. anomalous couplings

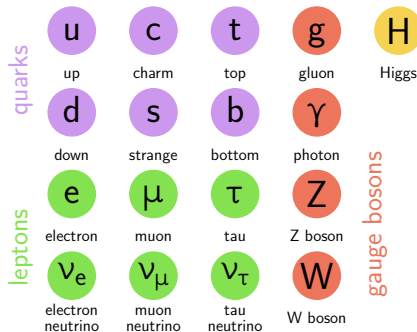


Requirements:

- ▶ Precise knowledge of the cross section and its dependences: scales, parton densities, couplings, masses, ...
- ▶ Quick calculation desirable!

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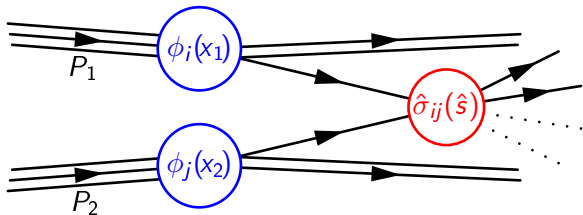


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Calculating Hadronic Cross Sections

Factorisation into PDFs and partonic cross section:

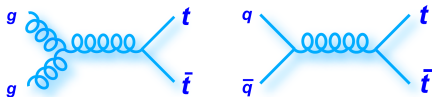


$$\sigma_{\text{had}} = \sum_{ij} \int_0^1 dx_1 \int_0^1 dx_2 \phi_{i,h_1}(x_1, \mu_F) \phi_{j,h_2}(x_2, \mu_F) \hat{\sigma}_{ij}(\hat{s}, \mu_F)$$

- ▶ $\phi(x, \mu_F)$ parton densities
- ▶ $\hat{\sigma}_{ij}(\hat{s}, \mu_F)$ partonic cross section
- ▶ $\hat{s} = x_1 x_2 s$ part. centre-of-mass energy
- ▶ $x_{1,2}$ fractional parton momentum
- ▶ μ_F factorisation scale

Top Pair Production

- ▶ Top anti-top pairs produced by **strong interaction**



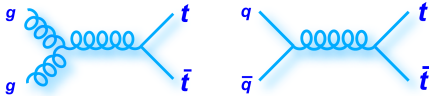
- ▶ Implemented in the **HATHOR** program
 - ▶ Fast calculation of hadronic cross sections for $t\bar{t}$ production:
[M. Aliev et al, Comput.Phys.Commun.182: 1034-1046, 2011]
 - ▶ Provides reference values in NNLO-QCD:

$$\hat{\sigma} = \underbrace{\alpha_s^2 \hat{\sigma}^{\text{LO}}}_{\text{leading order}} + \underbrace{\alpha_s^3 \hat{\sigma}^{\text{NLO}}}_{\text{next-to-leading order}} + \underbrace{\alpha_s^4 \hat{\sigma}^{\text{NNLO}}}_{\text{next-to-next-to-leading order}} + O(\alpha_s^5)$$

- ▶ Calculation of theoretical uncertainties (PDFs, scales, ...)

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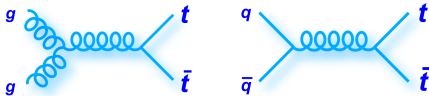
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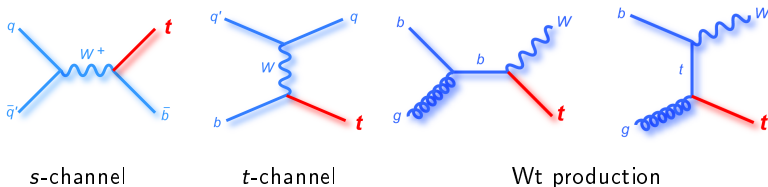
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Electroweak Top-Quark Production

Production of single top-quarks:

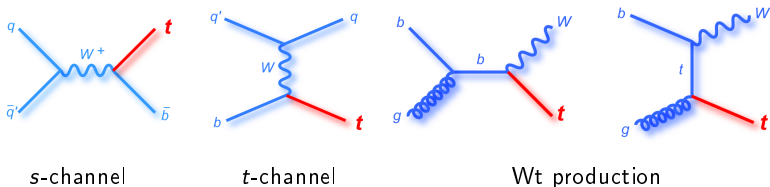


- ▶ NLO calculation exists in MCFM, not in HATHOR (Monte Carlo for FeMtobarn processes)
- ▶ Long running: $\mathcal{O}(1\text{h})$ per hadronic cross section
- ▶ Studying theoretical uncertainties: repeated execution necessary, requires $\mathcal{O}(\text{days})$

⇒ Desirable:
Faster calculation of total hadronic cross sections

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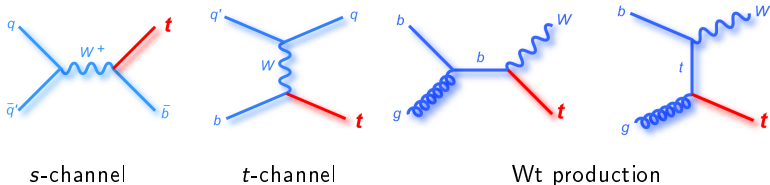


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Extraction of Partonic Cross Sections

$$\sigma_{\text{had}}(s) = \sum_{ij} \int dx_1 dx_2 \phi_{i,h_1}(x_1, \mu_F) \phi_{j,h_2}(x_2, \mu_F) \hat{\sigma}_{ij}(\underbrace{x_1 x_2 s}_{\hat{s}}, \mu_F)$$

For $t\bar{t}$ production:

- ▶ $\hat{\sigma}$ has closed-form expression (function of \hat{s})

For single-top quark production:

- ▶ Code for $\hat{\sigma}$ does not exist in closed-form
- ▶ $\hat{\sigma}(\hat{s})$ implemented in MCFM as follows:

$$\hat{\sigma} = \underbrace{\int_n d\sigma_{\text{LO}}}_{\text{easy, finite}} + \underbrace{\int_n d\sigma_{\text{virt.}} + \int_{n+1} d\sigma_{\text{real}} + \int_{n+1} d\sigma_{\text{fact.}}}_{\text{infrared/collinear divergences in each part}}$$

- ▶ Integration of real and virtual corrections on different phase spaces \Rightarrow Cancellation of divergences is not trivial

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- ▶ **Target:** Extract $\hat{\sigma} = \int d\hat{\sigma}(\hat{s})$
- ▶ **Idea:** Fix the partonic centre-of-mass energy \hat{s} in integration
- ▶ Minimal-invasive approach: replace PDFs with delta functions

$$\phi(x) \xrightarrow{\text{insert}} \delta(x - x_0) \approx \frac{1}{\sqrt{2\pi}w} \exp\left(-\frac{(x - x_0)^2}{2w^2}\right)$$

- ▶ MCFM calculates effectively $\hat{\sigma}$ at arbitrary position $x_0^2 s$
- ✓ Simple and quick implementation
- ✗ Inefficient in numerical integration (esp. at $w \rightarrow 0$)
- ✗ Systematic error due to finite width w of PDFs

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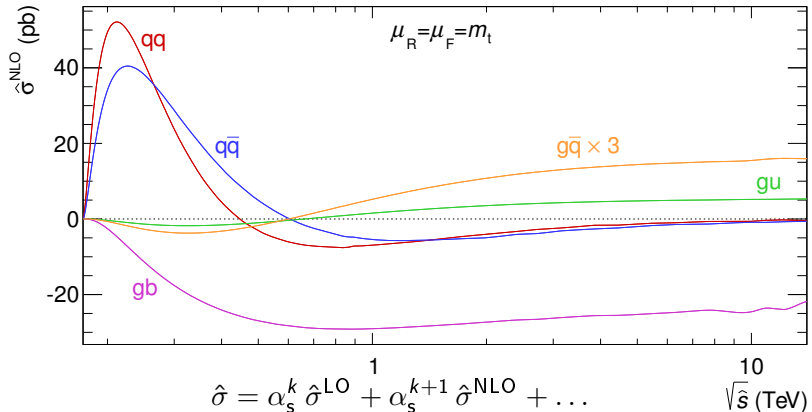
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Extraction of Partonic Cross Sections (I)

Result: NLO contributions in [single-top](#) t -channel



Timespan: [several weeks](#) for all processes!

Extraction of Partonic Cross Sections (II)

Alternative approach:

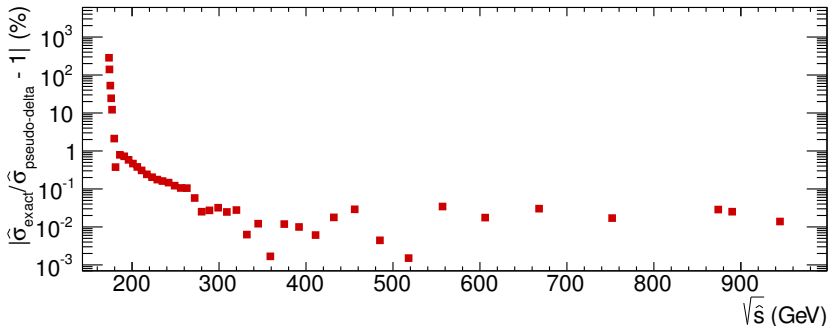
- ▶ Identification of partonic centre-of-mass energy in MCFM
- ▶ Insert exact delta functions:

$$\hat{\sigma}(\hat{s}) = \int \delta(x_1 - x_0) \delta(x_2 - x_0) d\hat{\sigma}(x_1 x_2 s) dx_1 dx_2$$

- ▶ Adjustments to $x_{1,2}$ directly in MCFM code
- ✓ Very efficient extraction von $\hat{\sigma}$
- ✓ No inherent, systematic error
- ✗ Invasive modifications to the code necessary

Comparison of Both Methods

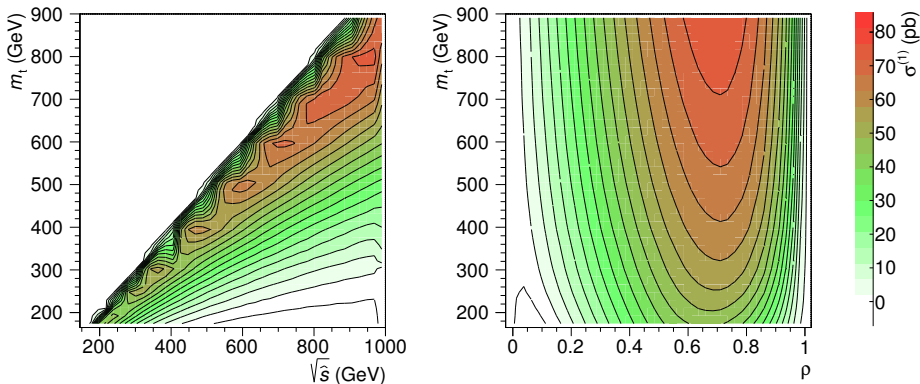
Relative difference of both methods in t -channel ($ub \rightarrow tq$)



- ▶ Successful consistency check, very good agreement
 - ▶ First method's issues at threshold already known
- ⇒ Using this (2nd) method:
- Extraction as function of \hat{s} and m_t possible

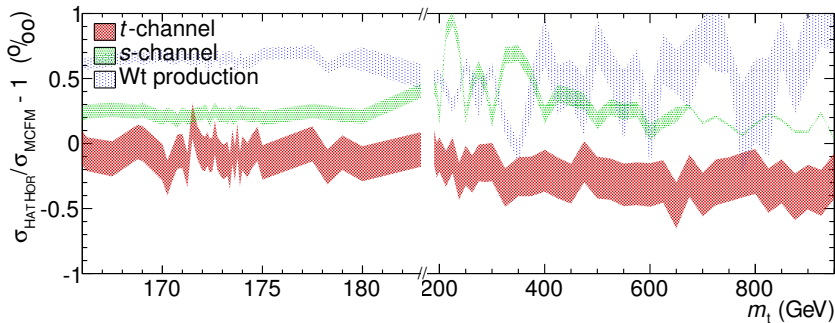
Partonic Cross Section

- ▶ Extraction of $\hat{\sigma}$ for many sampling points \hat{s} und m_t
- ▶ Implementation of $\hat{\sigma}$ in HATHOR
- ▶ Interpolation of $\hat{\sigma}(\hat{s}, m_t)$ on 2D grid necessary ($\rho = m_t^2 / \hat{s}$ helps)



- ▶ Integration of $\hat{\sigma}$ and PDFs in HATHOR $\Rightarrow \sigma_{\text{had}}$!

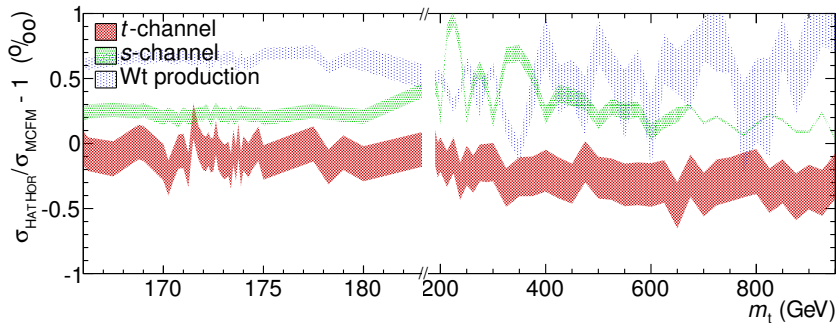
Consistency Check



- ▶ HATHOR: Convolution of extracted $\hat{\sigma}$ and PDFs
- ▶ Comparison of σ_{had} with complete MCFM calculation
- ▶ Precision of $\mathcal{O}(0,1\%)$ over large mass range

Time required: ▶ HATHOR: ~ 1 hour
▶ MCFM: ~ 1 week

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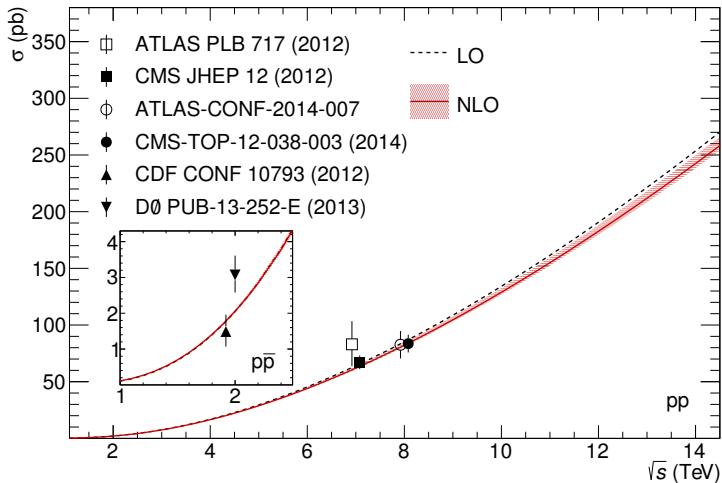
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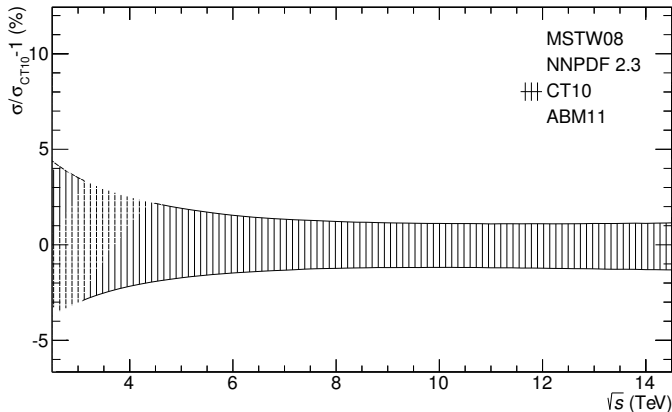
Summary

Hadronic Cross Section: t -channel



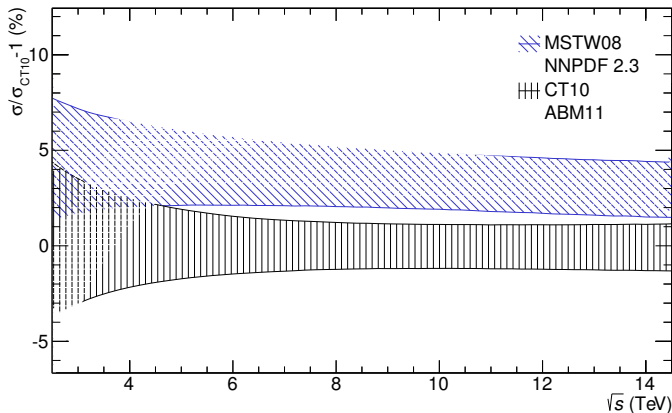
- Different orders of perturbation series in t -channel

Impact of PDFs in NLO



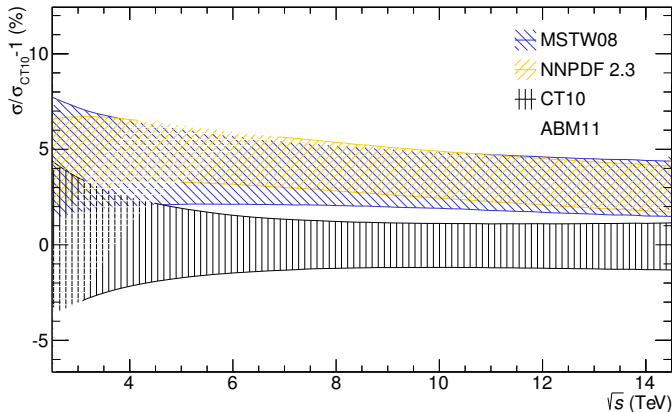
- ▶ PDF uncertainties in NLO (t -channel)
- ▶ Time consuming: ~ 100 error PDFs
- ▶ Errors from varying PDF fit parameters and α_s
- ▶ Error band does not always cover differences

Impact of PDFs in NLO



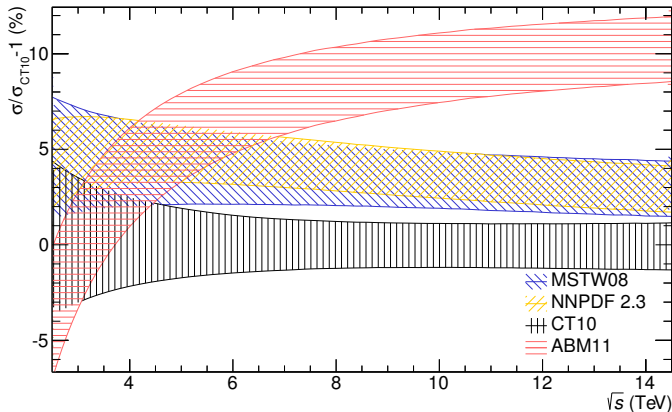
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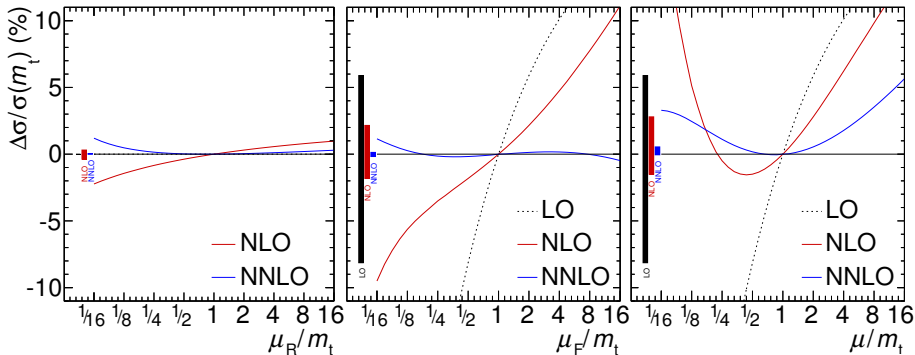
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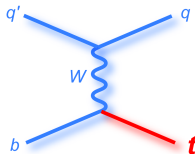


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Scale Dependence of the Cross Section



- Scale dependence in the *t*-channel (exact calculation by S. Mölbitz)
- Cross section normalized to natural scale ($\mu = m_t$)
- Dependences in higher orders considerably reduced



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Extraction of partonic cross sections:

- ✓ Implemented two methods
- ✓ Consistency checks using $t\bar{t}$ production
- ✓ Applied to production of single top-quarks in NLO

Implementation in HATHOR:

- ✓ Reached precision of $\leq \mathcal{O}(0.1\%)$
- ✓ Much better performance than e.g. MCFM
- ✓ Exact scale dependence in NLO und NNLO

Demonstrated abilities of the implementation:

- ✓ Study of systematic uncertainties (PDFs, scales, CKM, ...)
- ✓ Possibility to produce plots in $\mathcal{O}(1\text{ h})$
- ✓ Application to top-quark mass measurement

Release:

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 - Preprint online (arXiv:1406.4403), publication in progress

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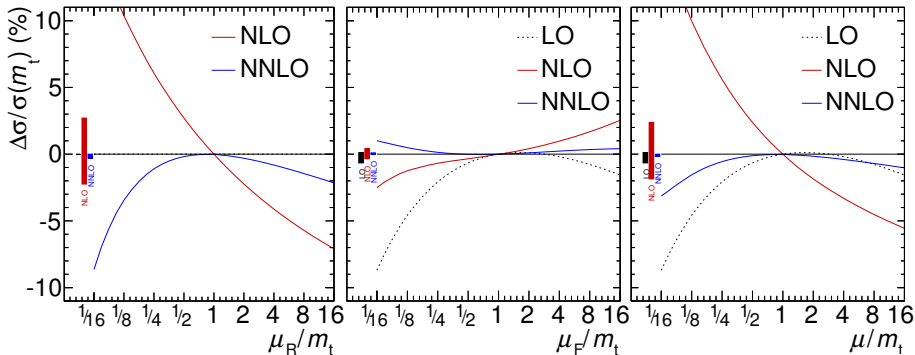
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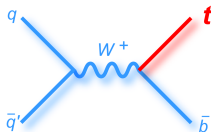
Thank you for your attention!

Backup

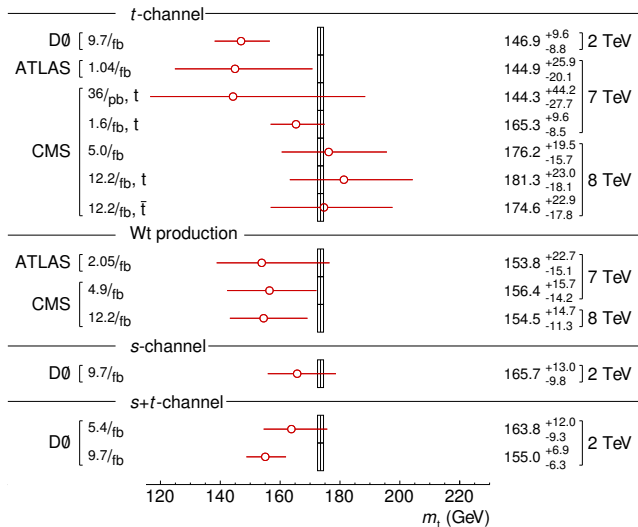
Scale Dependence of the Cross Section



- Scale dependence in the **s-channel** (exact calculation by S. Mölbitz)
- Cross section normalized to natural scale ($\mu = m_t$)
- Underlines necessity of higher-order calculations



Determination of the Top-Quark Mass



- Determining the top-quark mass from cross sections in **NLO**
- Correlations neglected
- Highest sensitivity in s-channel
- Lowest sensitivity in t-channel

Implementation of Scale Dependences

- ▶ Expand partonic cross section:

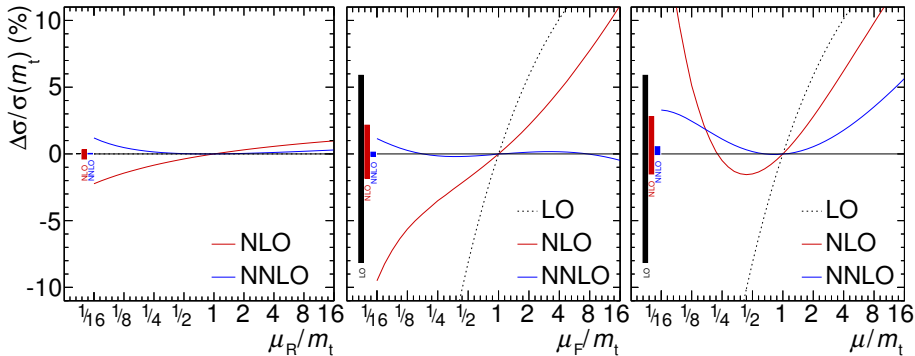
$$\hat{\sigma} = \alpha_s^k \hat{\sigma}^{\text{LO}} + \alpha_s^{k+1} \hat{\sigma}^{\text{NLO}} + \dots$$

- ▶ Extraction of $\hat{\sigma}^{\text{NLO}}$ at a scale $\mu_F = \mu_R = m_t$
- ▶ Extension to arbitrary scales:

$$\hat{\sigma}^{\text{NLO}} = \hat{\sigma}^{(10)} + \log\left(\frac{\mu_F^2}{m_t^2}\right) \hat{\sigma}^{(11)} + k\beta_0 \log\left(\frac{\mu_F^2}{\mu_R^2}\right) \hat{\sigma}^{\text{LO}}$$

- ▶ Implement $\hat{\sigma}^{(10)}$ using extracted cross sections
- ▶ Calculation of $\hat{\sigma}^{(11)}$ in master thesis of S. Mölbitz
- ▶ Allows variation of α_s and scales

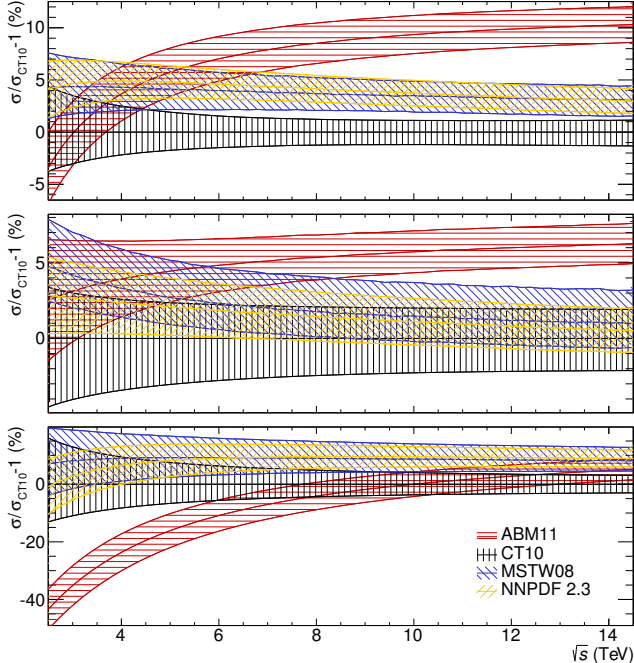
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- ▶ Scale dependence in *t*-channel
(exact calculation by S. Mölbitz)
- ▶ Very small μ_R dependence
- ▶ Dependences in NNLO considerably reduced

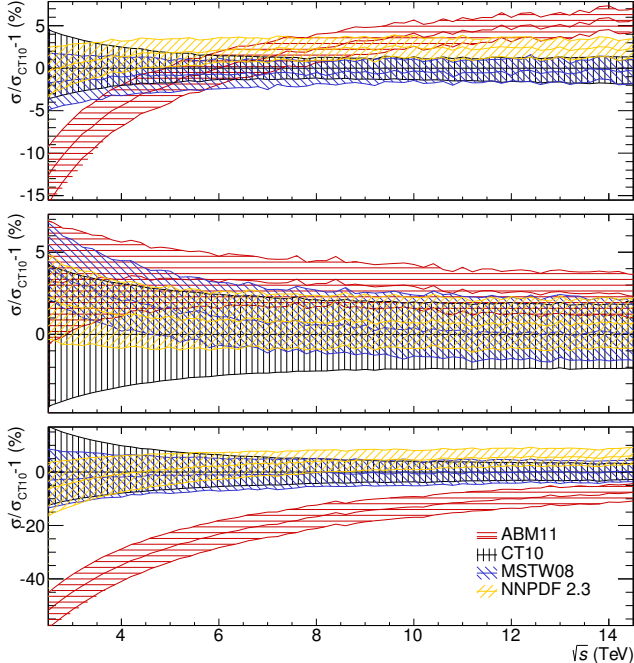
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- ▶ Error band covers differences better than in NLO

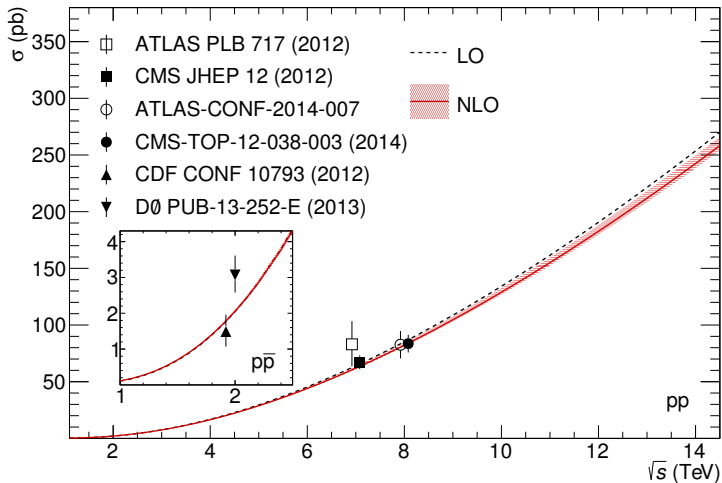


$$\begin{aligned}
 \sigma_{\text{had}}(S) = \sum_{ij} \int_0^1 \int_0^1 dx_1 dx_2 \left\{ \right. & \phi_{i,h_1}(x_1) \phi_{j,h_2}(x_2) \hat{\sigma}_{ij}(\hat{s}) \\
 & + \int_{x_1}^1 \frac{dz}{z} \phi_{i,h_1}\left(\frac{x_1}{z}\right) \phi_{j,h_2}(x_2) \hat{\sigma}_{ij}^{(z)}(\hat{s}) \\
 & \left. + \int_{x_2}^1 \frac{dz}{z} \phi_{i,h_1}(x_1) \phi_{j,h_2}\left(\frac{x_2}{z}\right) \hat{\sigma}_{ij}^{(z)}(\hat{s}) \right\}
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 &\quad \left. + \int_0^1 dz \left[\hat{\sigma}_{ij}^{(z)}((z x_1) x_2 S) + \hat{\sigma}_{ij}^{(z)}(x_1 (z x_2) S) \right] \right\} \\
 &= \sum_{ij} \int_0^1 \int_0^1 dx_1 dx_2 \phi_{i,h_1}(x_1) \phi_{j,h_2}(x_2) \left\{ \hat{\sigma}_{ij}(\hat{s}) + \right. \\
 &\quad \left. + 2 \int_0^1 dz \hat{\sigma}_{ij}^{(z)}((z x_1 x_2) S) \right\}
 \end{aligned}$$

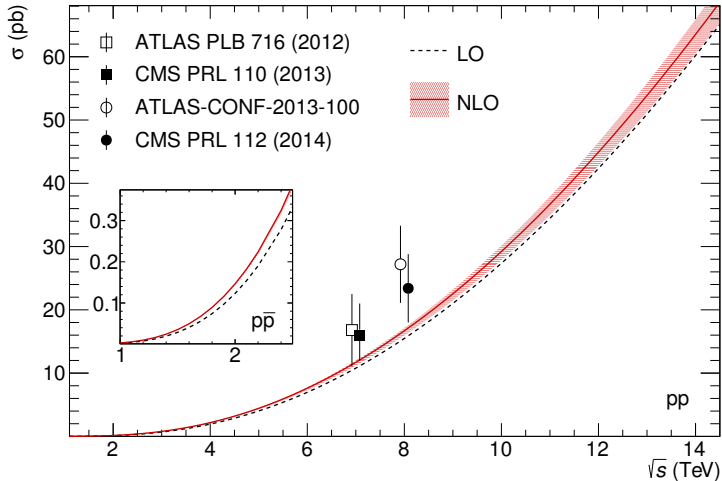
$$\hat{\sigma}_{ij}(\hat{s}) = \hat{\sigma}_{ij}(\hat{s}) + 2 \int_0^1 dz \, \hat{\sigma}_{ij}^{(z)}(z\hat{s})$$

Hadronischer Wirkungsquerschnitt: t -Kanal



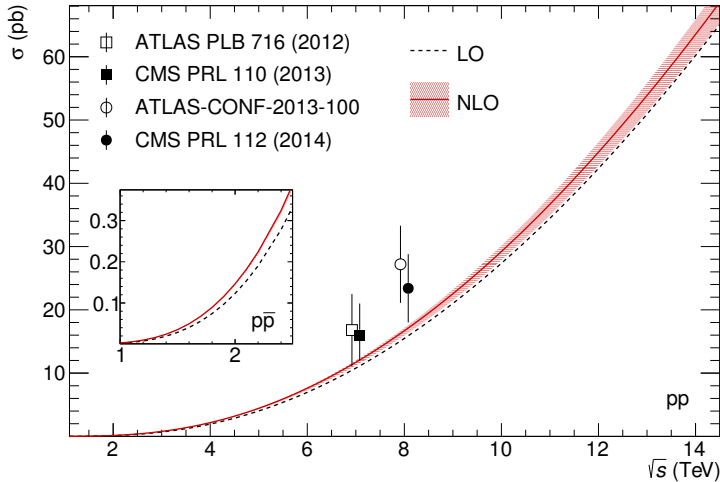
► Verschiedene Ordnungen der Störungsreihe im t -Kanal

Hadronischer Wirkungsquerschnitt: s-Kanal



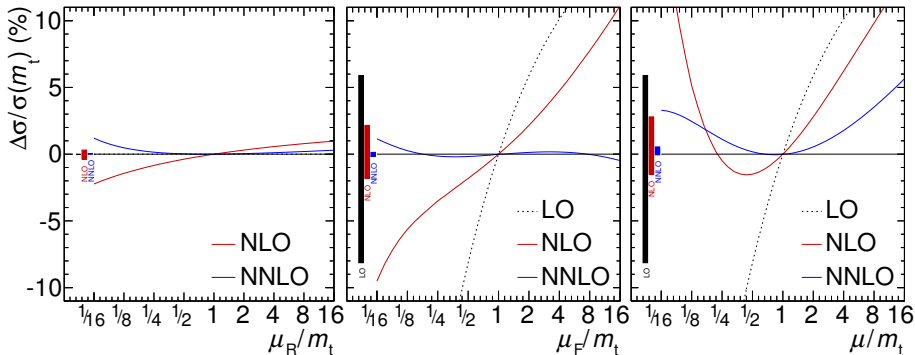
► Verschiedene Ordnungen der Störungsreihe im s-Kanal

Hadronischer Wirkungsquerschnitt: Wt -Produktion

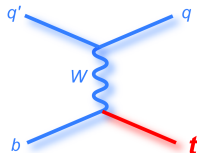


- Verschiedene Ordnungen der Störungsreihe für tW -Produktion

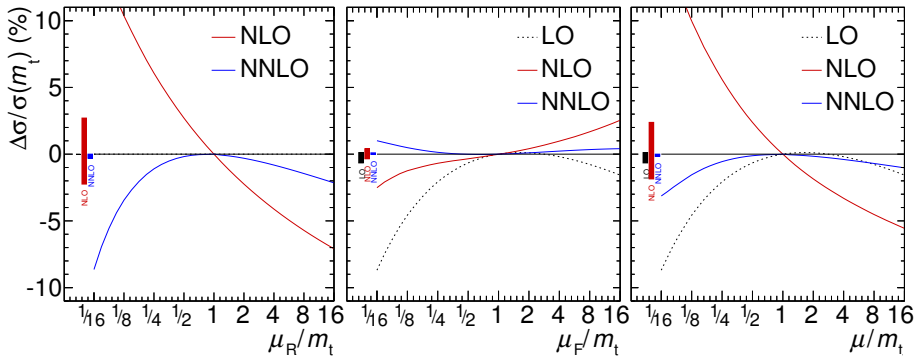
Skalenabhängigkeit des Wirkungsquerschnitts



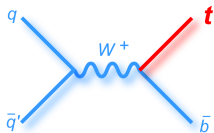
- Skalenabhängigkeiten im t -Kanal
(exakte Rechnung von S. Mölbitz)



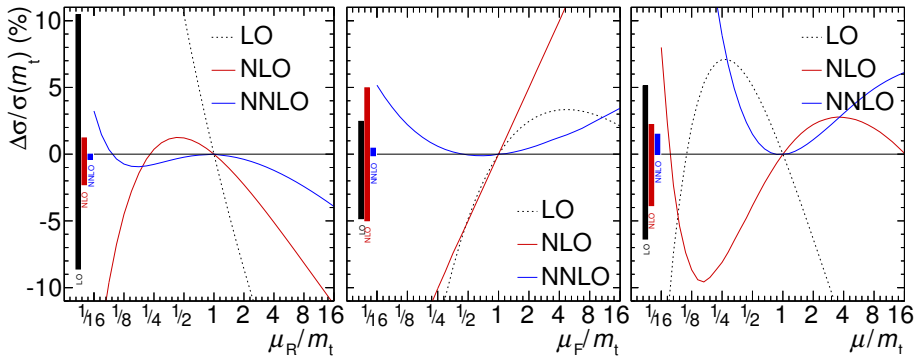
Skalenabhängigkeit des Wirkungsquerschnitts



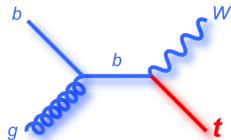
- Skalenabhängigkeiten im s-Kanal
(exakte Rechnung von S. Mölbitz)



Skalenabhängigkeit des Wirkungsquerschnitts



- Skalenabhängigkeiten für tW -Produktion (exakte Rechnung von S. Mölbitz)



Renormierungsskalenabhängigkeit

- ▶ Wirkungsquerschnitt in Potenzen von α_s :

$$\sigma = \alpha_s^k \sigma_0 + \alpha_s^{k+1} \sigma_1(\mu_R) + \alpha_s^{k+2} \sigma_2(\mu_R) + \mathcal{O}(\alpha_s^{k+3})$$

- ▶ β -Funktion der QCD:

$$\beta(\alpha_s) = \frac{\partial \alpha_s}{\partial \log \mu_R^2} = \mu_R^2 \frac{\partial \alpha_s}{\partial \mu_R^2} = -(4\pi)\beta_0 \alpha_s^2 - (4\pi)^2 \beta_1 \alpha_s^3 - \mathcal{O}(\alpha_s^4)$$

- ▶ Ableitung von σ nach μ_R :

$$\begin{aligned} \frac{\partial \sigma}{\partial \log \mu_R^2} &= k \alpha_s^{k-1} \beta(\alpha_s) \sigma_0 + (k+1) \alpha_s^k \beta(\alpha_s) \sigma_1(\mu_R) + \alpha_s^{k+1} \frac{\partial \sigma_1(\mu_R)}{\partial \log \mu_R^2} \\ &\quad + (k+2) \alpha_s^{k+1} \beta(\alpha_s) \sigma_2(\mu_R) + \alpha_s^{k+2} \frac{\partial \sigma_2(\mu_R)}{\partial \log \mu_R^2} + \mathcal{O}(\alpha_s^{k+3}) \end{aligned}$$

- ▶ σ soll nicht von μ_R abhängen:

$$0 = \alpha_s^{k+1} \left(-4\pi k \beta_0 \sigma_0 + \frac{\partial \sigma_1(\mu_R)}{\partial \log \mu_R^2} \right)$$

- ▶ σ soll nicht von μ_R abhängen:

$$0 = \alpha_s^{k+1} \left(-4\pi k \beta_0 \sigma_0 + \frac{\partial \sigma_1(\mu_R)}{\partial \log \mu_R^2} \right)$$

- ▶ Auflösen nach σ_1 :

$$\begin{aligned} \sigma_1(\mu_R) &= -4\pi k \beta_0 \sigma_0 \log \mu_R^2 + \tilde{\sigma}_1 \\ &= -4\pi k \beta_0 \sigma_0 \log \frac{\mu_R^2}{\mu_F^2} + \tilde{\sigma}_1(\mu_F^2) \end{aligned}$$

- ▶ μ_R Abhängigkeit der NLO-Korrektur ist proportional zu σ_0
- ▶ Integrationskonstante ist μ_R -unabhängiger Teil (aber noch $\sim \mu_F$)

- ▶ DGLAP-Gleichung ($2n_f + 1$ Dimensionen):

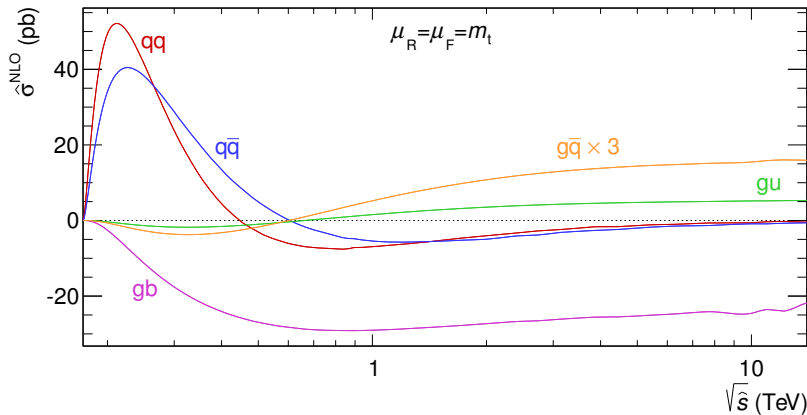
$$\frac{\partial}{\partial \log \mu_R^2} \begin{pmatrix} q_i(x, \mu_F) \\ g(x, \mu_F) \end{pmatrix} = 4\pi\alpha_s \sum_{q_i, \bar{q}_j} \begin{pmatrix} P_{q_i, q_j} & P_{q_i, g} \\ P_{g, q_j} & P_{gg} \end{pmatrix} \otimes \begin{pmatrix} q_i \\ g \end{pmatrix} (x)$$

- ▶ Entwickle P_{ij} in Potenzen von α_s
- ▶ Gleichungen müssen für alle PDF gelten
- ▶ Verlange für beide Skalen

$$\frac{\partial \sigma}{\partial \log \mu^2} = 0$$

Partonische Wirkungsquerschnitte

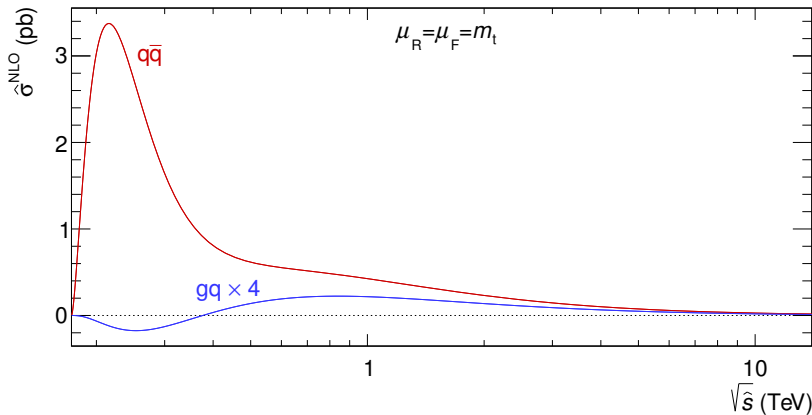
NLO-Beiträge im t -Kanal



$$\hat{\sigma} = \alpha_s^k \hat{\sigma}^{\text{LO}} + \alpha_s^{k+1} \hat{\sigma}^{\text{NLO}} + \dots$$

Partonische Wirkungsquerschnitte

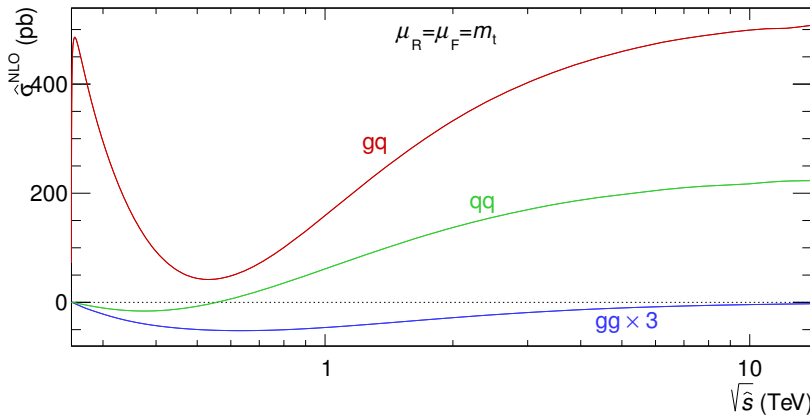
NLO-Beiträge im s -Kanal



$$\hat{\sigma} = \alpha_s^k \hat{\sigma}^{\text{LO}} + \alpha_s^{k+1} \hat{\sigma}^{\text{NLO}} + \dots$$

Partonische Wirkungsquerschnitte

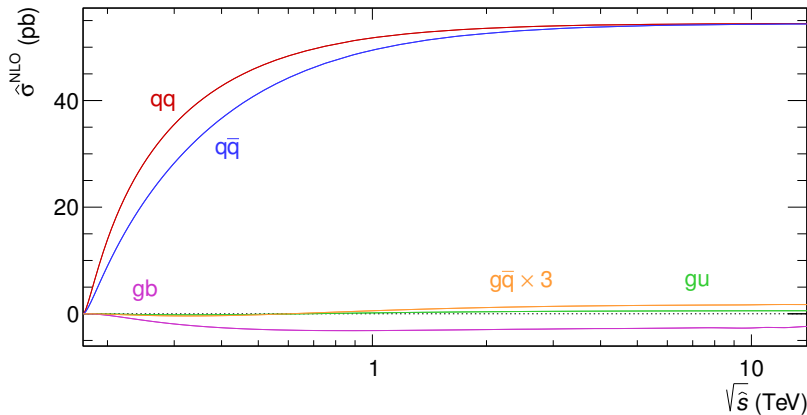
NLO-Beiträge zur Wt -Produktion



$$\hat{\sigma} = \alpha_s^k \hat{\sigma}^{\text{LO}} + \alpha_s^{k+1} \hat{\sigma}^{\text{NLO}} + \dots$$

Partonische Wirkungsquerschnitte

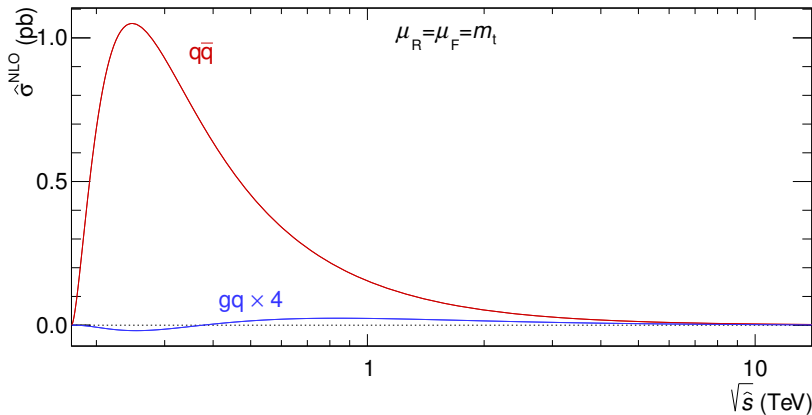
LO+NLO-Beiträge im t -Kanal



$$\hat{\sigma} = \alpha_s^k \hat{\sigma}^{\text{LO}} + \alpha_s^{k+1} \hat{\sigma}^{\text{NLO}} + \dots$$

Partonische Wirkungsquerschnitte

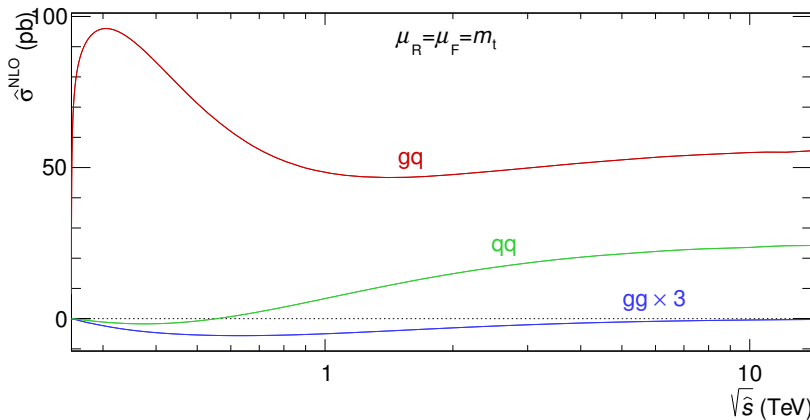
LO+NLO-Beiträge im s -Kanal



$$\hat{\sigma} = \alpha_s^k \hat{\sigma}^{\text{LO}} + \alpha_s^{k+1} \hat{\sigma}^{\text{NLO}} + \dots$$

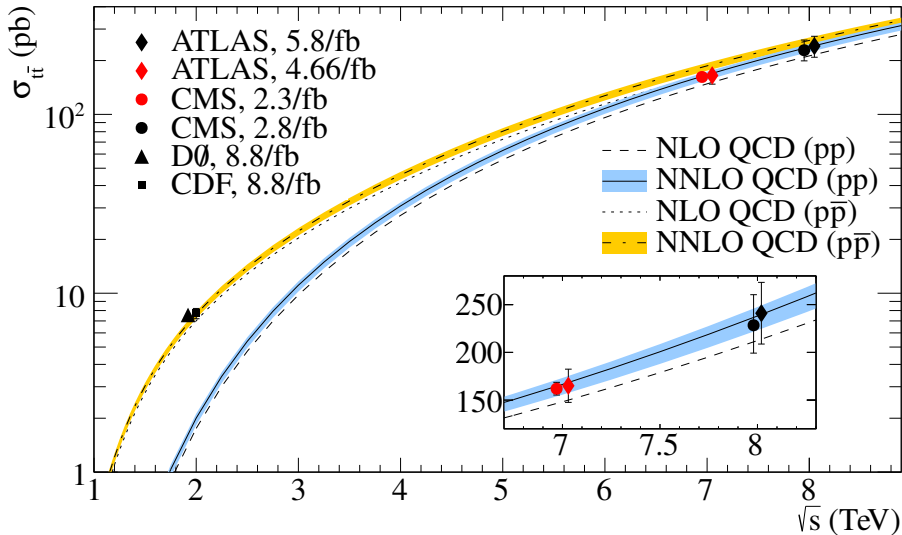
Partonische Wirkungsquerschnitte

LO+NLO-Beiträge zur Wt -Produktion

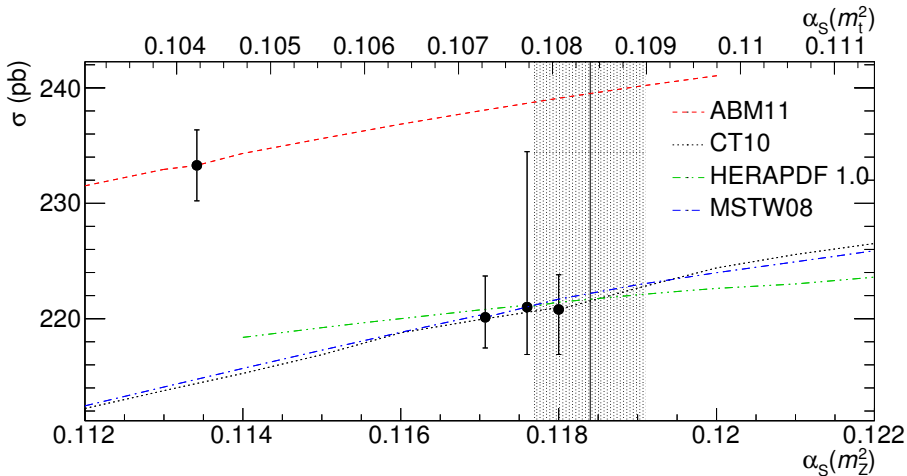


$$\hat{\sigma} = \alpha_s^k \hat{\sigma}^{\text{LO}} + \alpha_s^{k+1} \hat{\sigma}^{\text{NLO}} + \dots$$

HATHOR für Top-Paarherzeugung

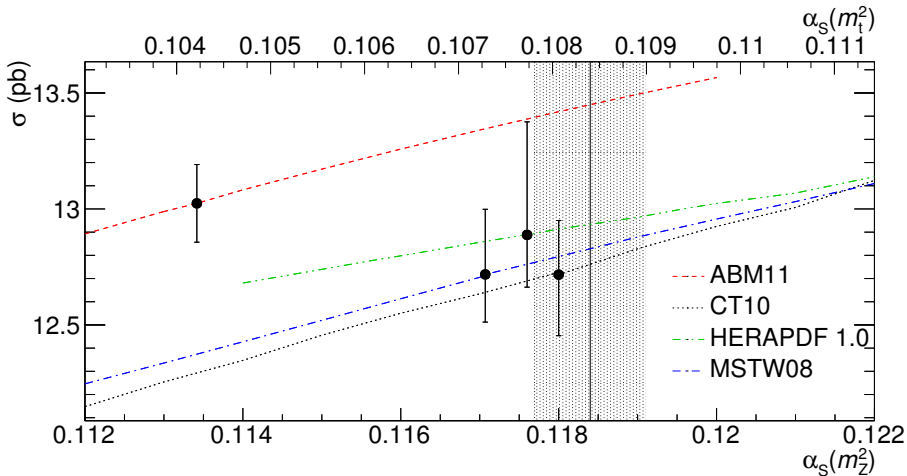


α_s -Abhängigkeit im t -Kanal



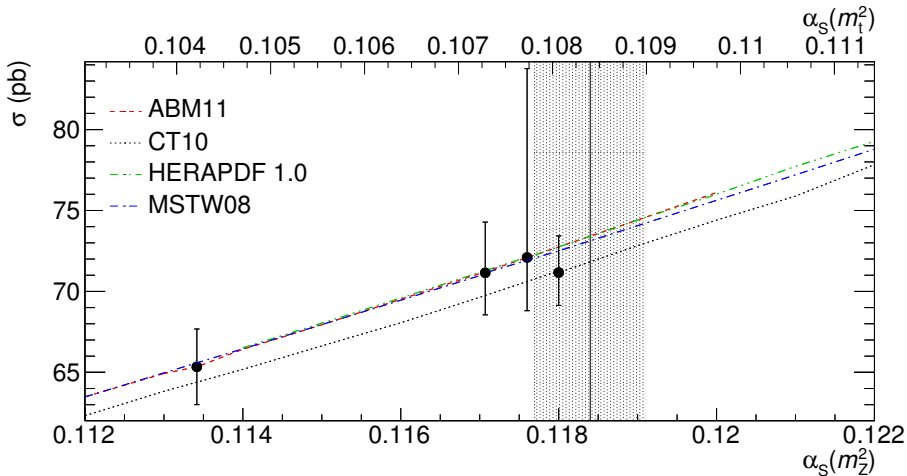
(in approx. NNLO, $\sqrt{s} = 14$ TeV)

α_s -Abhängigkeit im s-Kanal



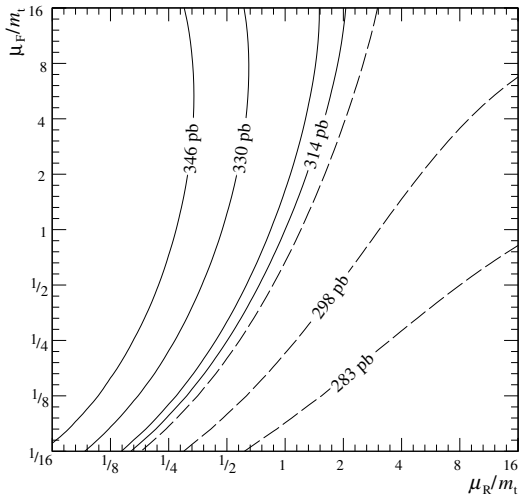
(in approx. NNLO, $\sqrt{s} = 14$ TeV)

α_s -Abhängigkeit der Wt -Produktion



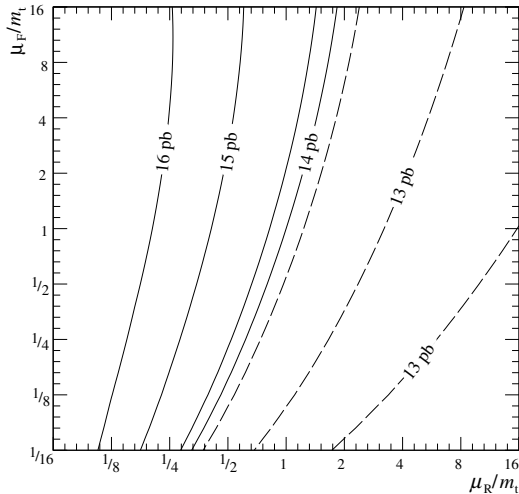
(in approx. NNLO, $\sqrt{s} = 14$ TeV)

Skalenabhängigkeit im t -Kanal



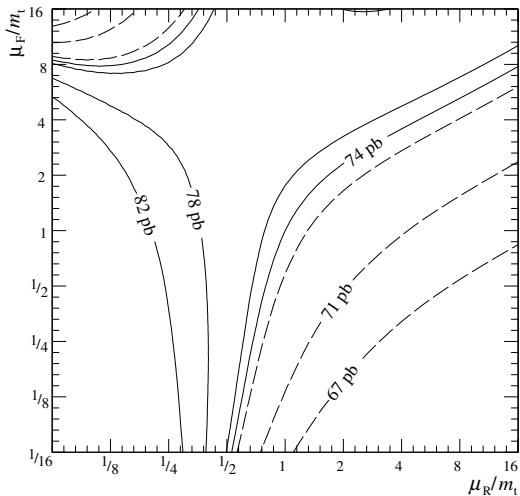
(in approx. NNLO, $\sqrt{s} = 14$ TeV)

Skalenabhängigkeit im s-Kanal



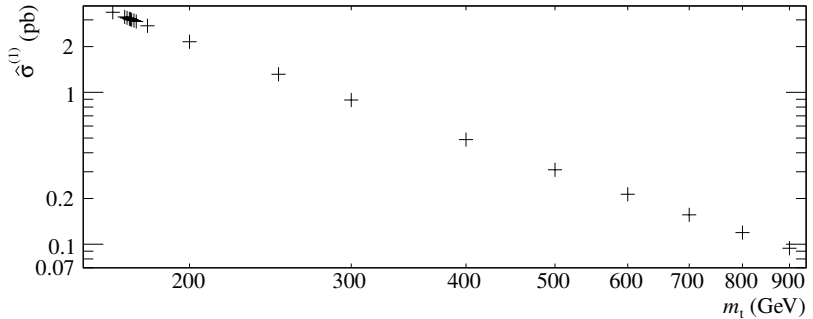
(in approx. NNLO, $\sqrt{s} = 14$ TeV)

Skalenabhängigkeit der Wt-Produktion

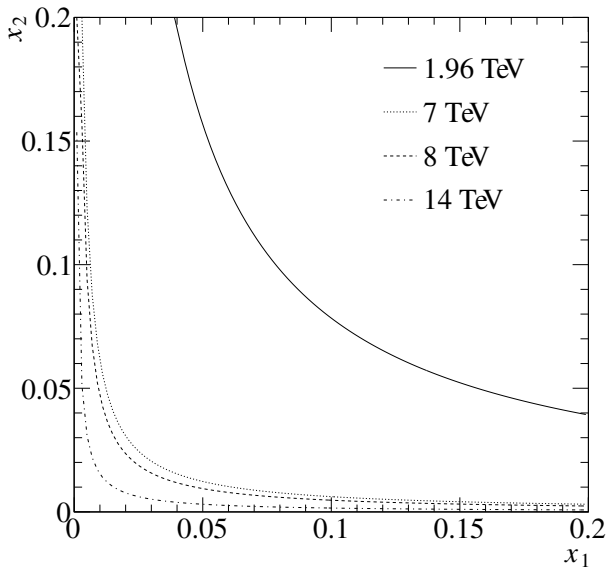


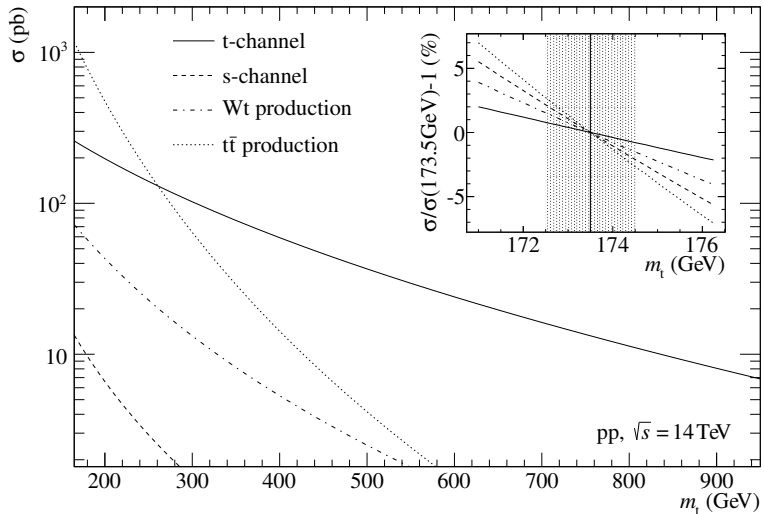
(in approx. NNLO, $\sqrt{s} = 14$ TeV)

Massenabhängigkeit (part. WQ)

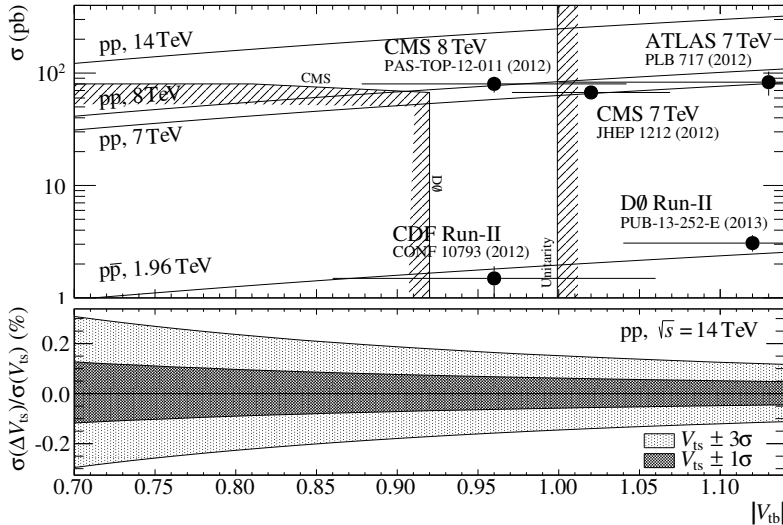


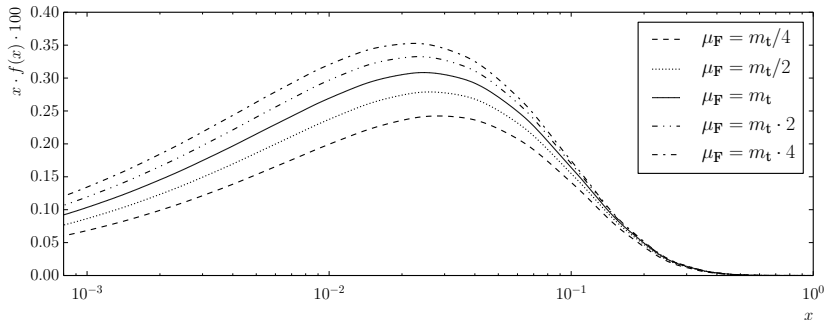
Zugänglicher Bereich in $x_{1,2}$





CKM-Matrixelementabhängigkeit





PDFs (Tevatron)

