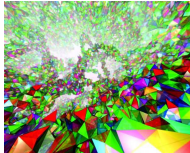


Loop Quantum Gravity – Part I: Introduction

Kristina Giesel

NORDITA

Astroteilchenphysik Schule 2009
Obertrubach-Bärnfels, 15.10.2009



Plan of the Talk

● Part I:

- Motivation for a theory of Quantum Gravity
- Candidate: Loop Quantum Gravity
- Classical starting point: General Relativity
- General Relativity & background independence
- Canonical quantisation of GR
- Quantum – Einstein – Equations and their interpretation

● Part II:

- Current research project within Loop Quantum Gravity
- Geometrical operators (volume, area & length)
- Solutions to Quantum – Einstein – Equations

● Summary & Conclusions

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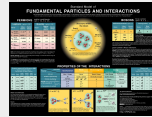
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Modern Theoretical Physics

- Two foundations: **QFT** and **GR**



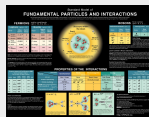
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- Describes fundamental interaction of electromagnetic, weak and strong force
- Description uses perturbative techniques (Exception: Lattice Gauge Theory)
- Infinities occur in perturbation series
- Renormalisation techniques can eliminate these infinities

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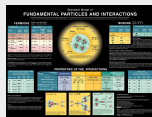
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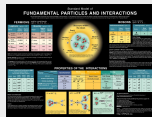
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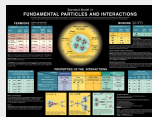
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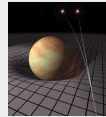
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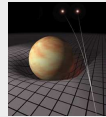
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- Describes fundamental interaction between matter & space time geometry
- Interaction is encoded in Einstein's equations
- Big Bang in cosmology, black holes \longrightarrow singularities
- Singularity theorems in General Relativity

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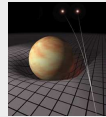
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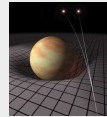
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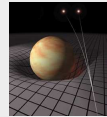
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Requirements for a Theory of Quantum Gravity:

- Adequate description of Quantum – Geometry & Quantum – Matter
- Principles of GR & QFT should be implemented consistently
- Background independence & canonical quantisation

Quest for a Theory of Quantum Gravity

Hope: Quantum Gravity is Fundamental Theory



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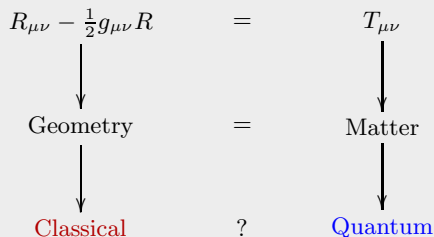


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Classical Einstein's – Equations

Interaction Between Space Time Geometry & Matter



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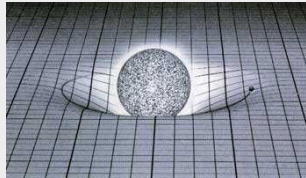
General Relativity

Basics of General Relativity

- **General Covariance:** Physics does not depend on coordinates

Einstein's Equations

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = T_{\mu\nu}$$



- Metric $g_{\mu\nu}$ encodes space time geometry
- **NO** preferred background metric, metric becomes dynamical object
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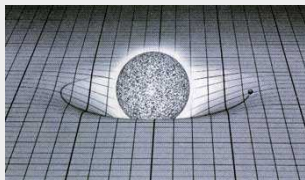
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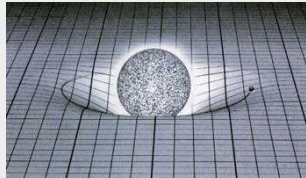
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Metric

Metric encodes geometry of spacetime

- In particle physics we assume Minkowski spacetime

$$\eta_{\mu\nu} = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

- Friedmann - Robertson - Walker - Metric (FRW)

$$g_{\mu\nu} = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & a(t) & 0 & 0 \\ 0 & 0 & a(t) & 0 \\ 0 & 0 & 0 & a(t) \end{pmatrix}$$

- Line element: $ds^2 = g_{\mu\nu} dx^\mu dx^\nu$
- $ds_M^2 = -dt^2 + dx^2 + dy^2 + dz^2$, $ds_{FRW}^2 = -dt^2 + a(t)(dx^2 + dy^2 + dz^2)$

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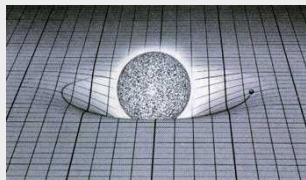
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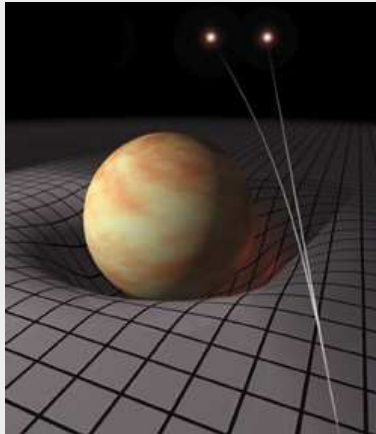
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Background Independence

Geometry and thus metric becomes dynamical object



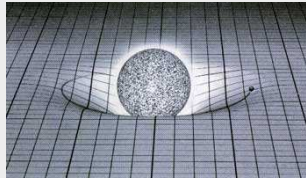
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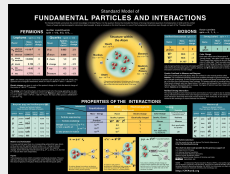
Quantum Field Theory

Basics of Quantum Field Theory

- QFT combines classical Field Theories & Quantum Mechanics

Canonical Quantisation in QM

$$\{q_j, p^k\} = \delta_j^k \longrightarrow [\hat{q}_j, \hat{p}^k] = i\hbar\delta_j^k$$



- preferred background metric: Minkowski
- Wightman Axioms rely on high symmetry of Minkowski-metric
- Background – dependent theory, perturbative theory
- Quantum theory: Canonical quantisation (Path integrals)

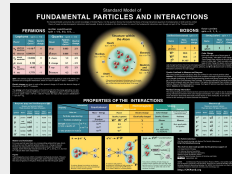
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Canonical Quantisation

Canonical Quantisation: Quantum Mechanics

- Elementary phase space variables q_i, p^j satisfy Poisson algebra $\{q_i, p^j\} = \delta_i^j$
- Dynamics described by Hamiltonian $H(p, q)$, example harmonic oscill.
 $H(p, q) = \frac{p^2}{2m} + \frac{1}{2}m\omega q^2$
- Construct a corresponding quantum algebra which respects the usual canonical quantisation rule $[\hat{q}, \hat{p}] = i\hbar$
- Look for possible representations of this quantum algebra
- Operators $\hat{q}\psi(x) := x\psi(x)$, $\hat{p}\psi(x) := -i\hbar \frac{d}{dx}\psi(x)$
- Stone – von Neumann – theorem for QM: Schrödinger representation unique

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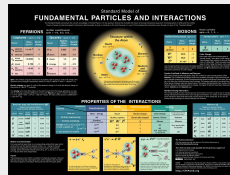
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$$\begin{aligned} \{\varphi(x), \pi(x')\} &= \delta(x, x') \longrightarrow \\ [\hat{\varphi}(x), \hat{\pi}(x')] &= i\hbar\delta(x, x') \end{aligned}$$



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- Smear $\phi(x)$ and $\pi(x')$ with test functions
- $\phi(f) := \int d^3x \phi(x) f(x)$, $\pi(f) = \int d^3x \pi(x) f(x)$
- Construct a corresponding quantum algebra and look for representations
- No Stone – von Neumann – theorem in this case
- Usually works with Fock space: $a(f) = \phi(f) + i\pi(f)$, $a(f)|0\rangle = 0 \dots$

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Canonical Quantisation

Canonical Quantisation: Quantum Field Theory

- Elementary phase space variables $\phi(x), \pi(x')$ satisfy $\{\phi(x), \pi(x')\} = \delta(x, x')$
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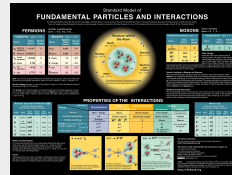
Basics of Quantum Field Theory

- QFT combines classical Field Theories & Quantum Mechanics

Canonical Quantisation in QM

$$\{q_j, p^k\} = \delta_j^k \longrightarrow [\hat{q}_j, \hat{p}^k] = i\hbar\delta_j^k$$

$$\{\varphi(x), \pi(x')\} = \delta(x, x') \longrightarrow [\hat{\varphi}(x), \hat{\pi}(x')] = i\hbar\delta(x, x')$$



- **preferred** background metric: Minkowski
- Wightman Axioms rely on high symmetry of Minkowski-metric
- Background – dependent theory, perturbative theory
- Quantum theory: Canonical quantisation (Path integrals)

Quantum Field Theory

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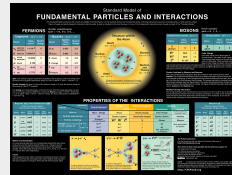
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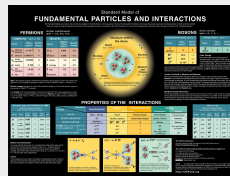
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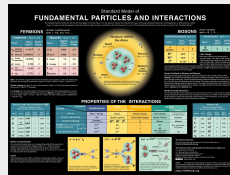
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Request for Theory of Quantum Gravity

Are we only bothered by occurrence of infinities?

- Anisotropies in cosmic microwave background
- Cosmological constant problem
- Black holes, Hawking radiation

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Loop Quantum Gravity

Loop Quantum Gravity: Candidate for a Quantum Theory of GR



Loop Quantum Gravity

1. Step:

- Derive Hamiltonian formulation of General Relativity

2. Step:

- Apply canonical quantisation to Hamiltonian form of General Relativity

3. Step:

- Quantum analogue of classical Einstein's Equations: Quantum – Einstein – Equations

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Classical Starting Point of GR

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$$S = \int_{\mathcal{M}} d^4x \sqrt{|\det(g)|} R$$

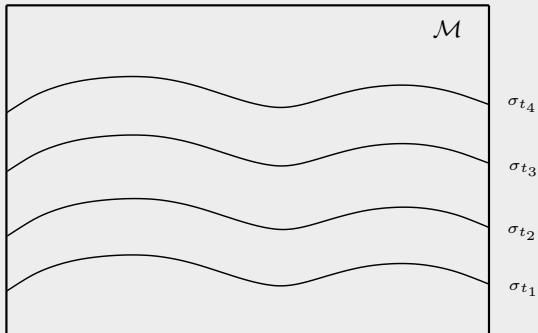
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Hamiltonian Formulation (ADM – formalisms)

- Split into space & time: $\int_{\mathcal{M}} \rightarrow \int_{\mathbb{R}} \int_{\sigma}$
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Foliation of Space Time

(3+1) Split into Space & Time



Classical Formulation of GR

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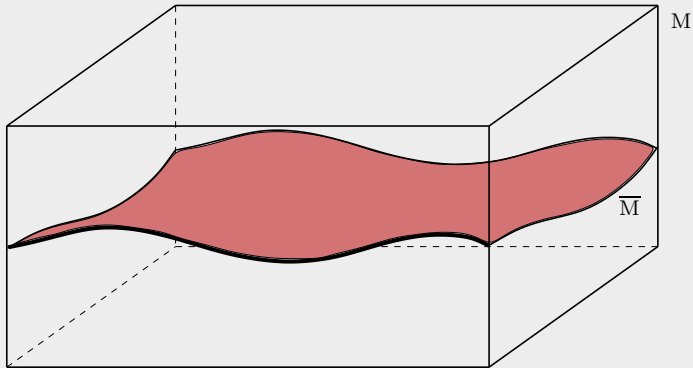
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Classical Formulation of GR

Constraint Hypersurface



Langrangian and Hamiltonian Picture

Langrangian Framework

- Einstein equations: $R_{\mu\nu}(g) - \frac{1}{2}g_{\mu\nu}R(g) = 0$
- $g_{\mu\nu}$ has 10 independent components \longrightarrow 10 equations
- Only six of them are second order differential equations

Hamiltonian Framework

- Hamiltonian Equations: $\dot{q}_{ab} = \{q_{ab}, H_{\text{can}}\}$, $\dot{p}^{ab} = \{p^{ab}, H_{\text{can}}\}$
- 12 first order differential equations \longrightarrow 6 second order for q_{ab}
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Constraints of General Relativity

Explicit Form in ADM Variables for Vacuum Case:

- Diffeomorphism constraint

$$D_a(q, p) = -2q_{ac}p^b_{;b}$$

- Hamiltonian constraint

$$H(q, p) = \frac{1}{\sqrt{\det(q)}} \left(q_{ac}q_{bd} - \frac{1}{2}q_{ab}q_{cd} \right) p^{ab}p^{cd} - \sqrt{\det(q)}R(q)$$

Classical Formulation of GR

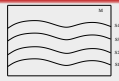
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- Split into space & time: $\int_{\mathcal{M}} \rightarrow \int_{\mathbb{R}} \int_{\sigma}$



- Elementary variables: $A_a^j(x)$, $E_k^b(x)$ $\{A_a^j(x), E_j^b(k)\} = \delta^3(x, y) \delta_a^b \delta_k^j$
- Constraints: $G_j(A, E) = 0$, $D_a(A, E) = 0$, $H(A, E) = 0$
- Holonomy – flux – algebra $\{A(e), E(S)\}$
- Starting point for quantisation for LQG

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Explicit Form in Ashtekar – Variables for Vacuum Case:

- Gauß constraint

$$G_j(A, E) = D_a E_j^a$$

- Diffeomorphism constraint

$$D_a(A, E) = F_{ab}^j E_j^b \quad F_{ab}^j \text{ curvature of } A_a^j$$

- Hamiltonian constraint

$$H(A, E) = \frac{1}{\sqrt{\det(E)}} \left(\epsilon_i^{jk} F_{ab}^i E_j^a E_k^b \right)$$

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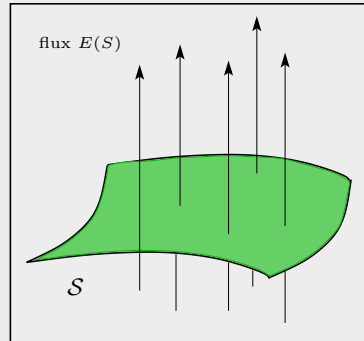
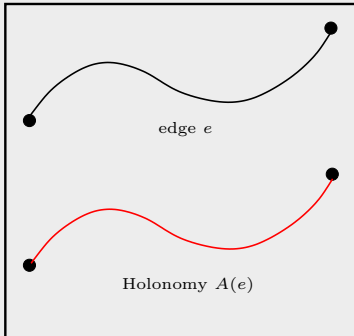
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Elementary phase space variables of LQG

Holonomies $A(e)$ and Fluxes $E(S)$



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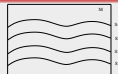
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Canonical Quantisation for Constrained Systems

Option 1: Dirac's Programme for Constraint Quantisation

- Quantise kinematical theory $\{q_{ab}(x), p^{cd}(x')\} = \delta^3(x, x') \delta_a^c \delta_d^b$
- Solve constraints in Quantum theory
one gets a kinematical Hilbert space \mathcal{H}_{kin}
- Physical states are then $\hat{C}\psi_{\text{phys}} = 0$
- Physical states live in physical Hilbert space $\mathcal{H}_{\text{phys}}$
- **Positive:** Kinematical algebra simple
- **Negative:** Implementation of operators \hat{C} might be hard

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- Solve constraint classically, Poisson algebra of observables $\{O_A(x), O_E(x')\}$
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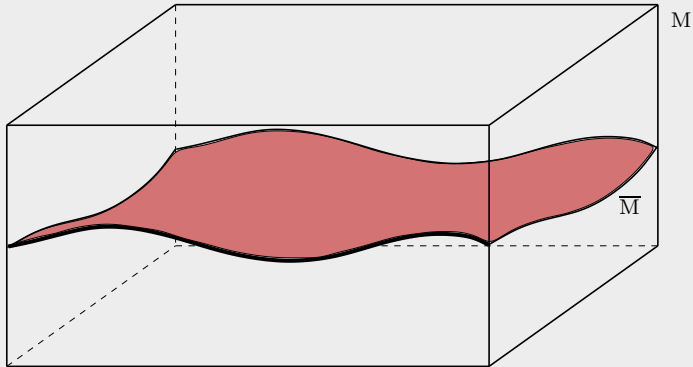
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Two Options for Constrained Systems

Option 1 & 2: Quantising M Or Only \overline{M}



Quantisation for Loop Quantum Gravity

Option 1 or Option 2?

- It does not mean that one of the options is preferred
- Different strategies to quantise systems with constraints
- Even a combination of both strategies might be useful:
Parts of the constraints are solved classically and the remaining ones in quantum theory

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Canonical Quantisation for LQG

Option 2: Dirac Quantisation for Loop Quantum Gravity

- Start with kinematical holonomy – flux – algebra $\{A(e), E(S)\}$
- One gets kinematical Hilbert space $\mathcal{H}_{\text{kin}} = L_2(\overline{\mathcal{A}}, d\mu_{\text{AL}})$
 [Ashtekar, Isham, Lewandowski, Rovelli, Smolin '90]
- LOST – Theorem [Lewandowski, Okolow, Sahlmann, Thiemann '05, Fleischhack '05]
- $\hat{A}(e)$ multiplication $\hat{E}(S)$ derivation operator
- In \mathcal{H}_{kin} implementation of the operators \hat{G}, \hat{D} und \hat{H}

Quantum – Einstein – Equations

$$\hat{G}_I \psi_{\text{phys}}(A) = 0, \quad \hat{D}_a \psi_{\text{phys}}(A) = 0 \quad \hat{H} \psi_{\text{phys}}(A) = 0$$

- Spin network functions $\psi(A)$ are ONB in \mathcal{H}_{kin}

Canonical Quantisation for LQG

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Hilbert Space

Recall from Quantum Mechanics

- Hilbert space $L_2(\mathbb{R}, dx)$ all functions for which

$$\int_{\mathbb{R}} d^3x |\psi(x)|^2 < \infty$$

- Here $L_2(\overline{A}, d\mu_{AL})$ all functions for which

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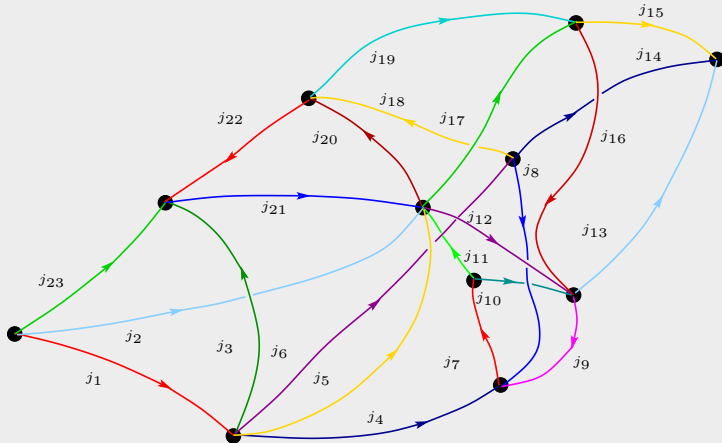
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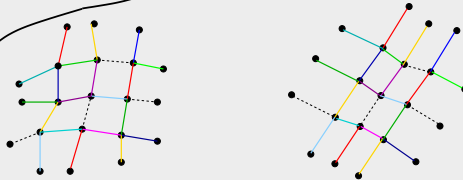
Basis of \mathcal{H}_{kin}

Spin network functions [Ashtekar, Isham, Lewandowski, Rovelli, Smolin '96]



LQG involves all embedded graphs

Embedded Graphs



Summary Part I

Summary

- LQG tries to combine background independence & canonical quantisation rigorously
- Starts with Hamiltonian formulation of General Relativity
 - ADM variables are not suitable for quantisation
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