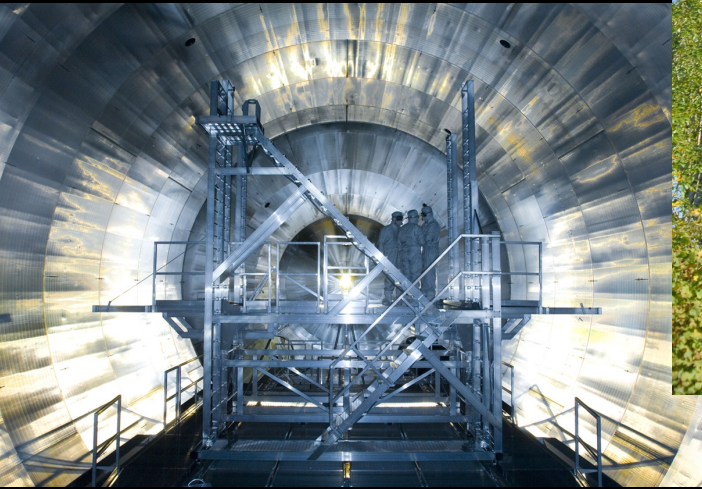


# keV sterile neutrino detection with KATRIN

Nicholas Steinbrink  
Institut für Kernphysik  
Uni Münster



Schule für Astroteilchenphysik 2013  
Obertrubach-Bärnfels



# Is dark matter neither hot nor cold, but warm (WDM)?

Cold dark matter (CDM) predicts too much substructure  
e.g. Anderhalden et al. 2013, <http://arxiv.org/abs/1212.2967>

**CDM**

**WDM**

Possible DM candidate: **sterile neutrino with  $\sim$  keV mass**

# Non trivial $\nu$ mixing

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

**Active**  $\nu$  flavours

Unitary mixing matrix

Mass states

# 3 active $\nu$ plus a sterile $\nu$

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \\ \nu_s \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & 0 \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & 0 \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \end{pmatrix}$$

New "flavour" state  $\nu_s$ :  
 Right handed  
 No coupling to W and Z

3x3 sub-matrix  
 still unitary

4<sup>th</sup> mass state  
 Identical to  $\nu_s$

# Non vanishing mixing between active and sterile $\nu$

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \\ \nu_s \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \end{pmatrix}$$

3x3 sub-matrix  
not unitary any longer

=> oscillation between  $\nu_{e/\mu/\tau}$  and  $\nu_s$  possible

=> new mass state  $\nu_4$  mixes with active flavour states

# Non vanishing mixing between active and sterile $\nu$

**Attention!** Some call  $\nu_s$  and some call  $\nu_4$  the "sterile neutrino"  
I will use both versions in the talk

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \\ \nu_s \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \end{pmatrix}$$

3x3 sub-matrix  
not unitary any longer

=> oscillation between  $\nu_{e/\mu/\tau}$  and  $\nu_s$  possible

=> new mass state  $\nu_4$  mixes with active flavour states

# Sterile $\nu$ mass scales

- **eV:**
  - favored for explaining reactor  $\nu$  / LSND / MiniBoone anomalies
- **keV:**
  - favored as dark matter candidate
- **above:**
  - leptogenesis, GUT...

Example of model with 3 sterile  $\nu$ :

*Canetti, Drewes, Frossard, Shaposhnikov Phys. Rev. D 87, 093006 (2013), arXiv: 1208.4607*

# Sterile $\nu$ mass scales

- **eV:**
  - favored for explaining reactor  $\nu$  / LSND / MiniBoone anomalies
- **keV:**
  - favored as dark matter candidate
- **above:**
  - leptogenesis, GUT...

Example of model with 3 sterile  $\nu$ :

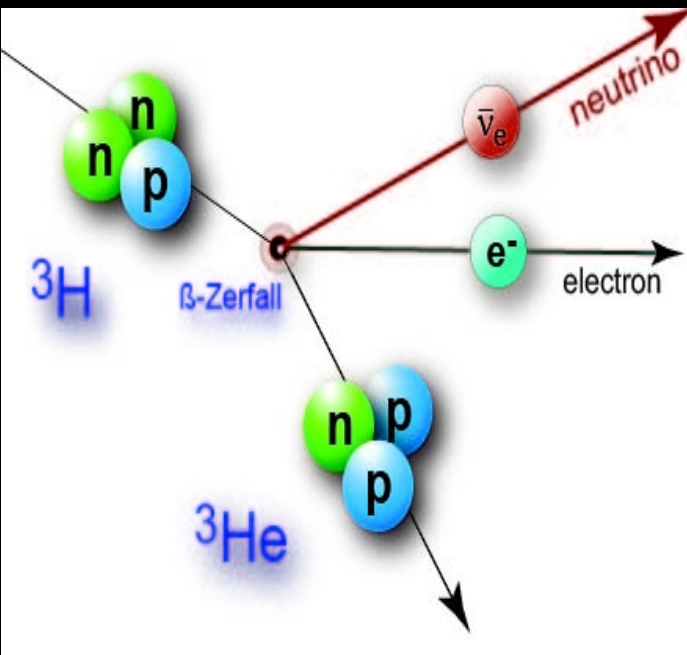
*Canetti, Drewes, Frossard, Shaposhnikov Phys. Rev. D 87, 093006 (2013), arXiv: 1208.4607*



# The tritium $\beta$ spectrum

$$\frac{dR}{dE}(E) = N \frac{G_F^2}{2\pi^3 \hbar^7 c^5} \cos^2(\theta_C) |M|^2 F(E, Z') \cdot$$

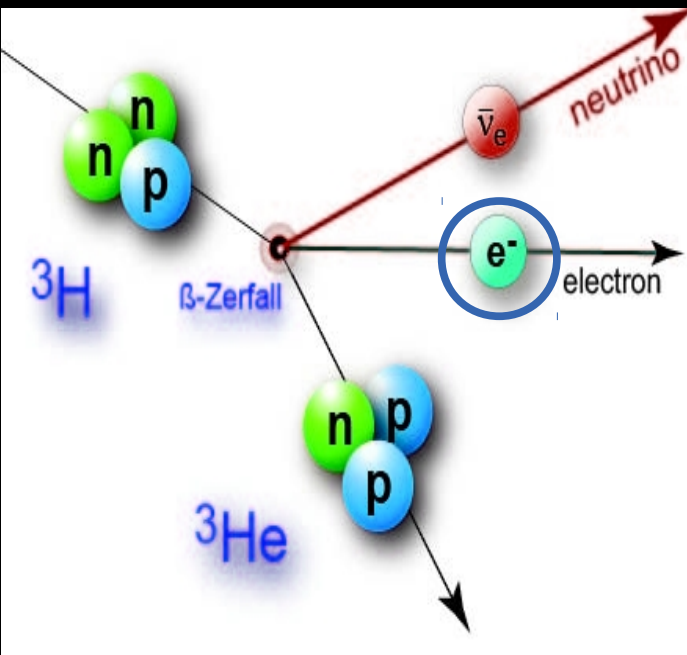
$$p \cdot (E + m_e c^2) \cdot (E_0 - E) \cdot \sum_i^3 U_{ei} \sqrt{(E_0 - E)^2 - m(\nu_i)^2 c^4}$$



# The tritium $\beta$ spectrum

$$\frac{dR}{dE}(E) = N \frac{G_F^2}{2\pi^3 \hbar^7 c^5} \cos^2(\theta_C) |M|^2 F(E, Z') \cdot$$

$$p \cdot (E + m_e c^2) \cdot (E_0 - E) \cdot \sum_i^3 U_{ei} \sqrt{(E_0 - E)^2 - m(\nu_i)^2 c^4}$$



# The tritium $\beta$ spectrum

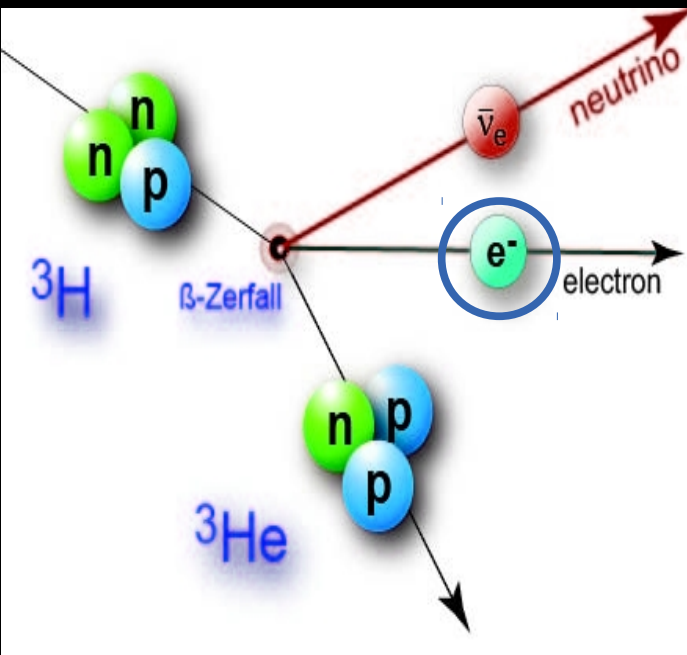
$$\frac{dR}{dE}(E) = N \frac{G_F^2}{2\pi^3 \hbar^7 c^5} \cos^2(\theta_C) |M|^2 F(E, Z') \cdot$$

$$p \cdot (E + m_e c^2) \cdot (E_0 - E) \cdot \sum_i^3 U_{ei} \sqrt{(E_0 - E)^2 - m(\nu_i)^2 c^4}$$

Def.: "electron neutrino mass"

$$m(\nu_e)^2 = \sum_i^3 |U_{ei}^2| m_i^2$$

$$\approx \sqrt{(E_0 - E)^2 - m(\nu_e)^2 c^4}$$



# The tritium $\beta$ spectrum

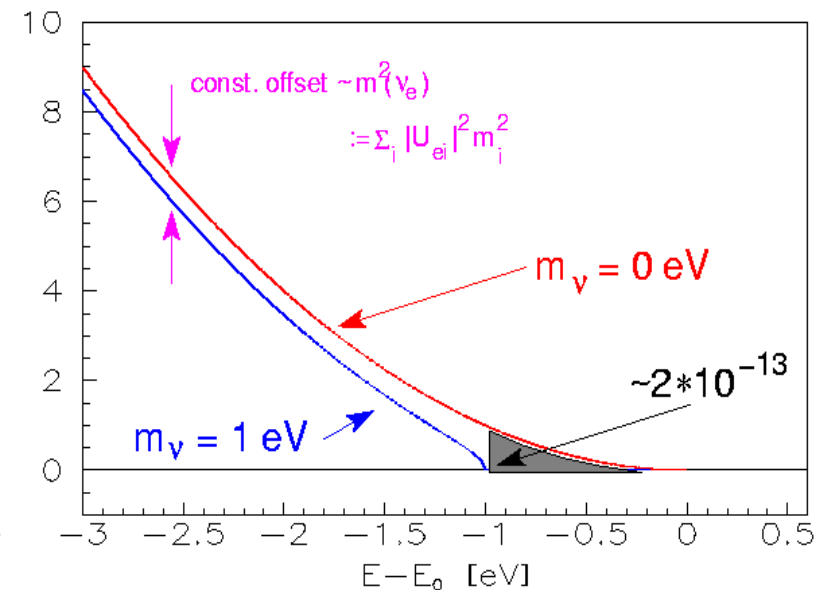
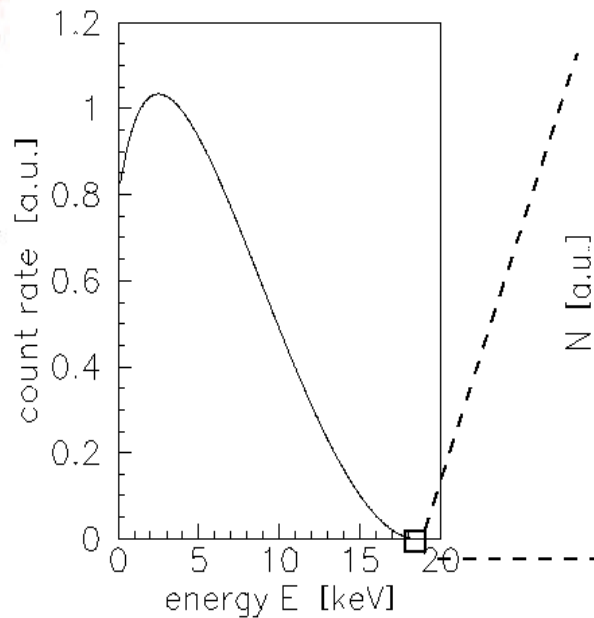
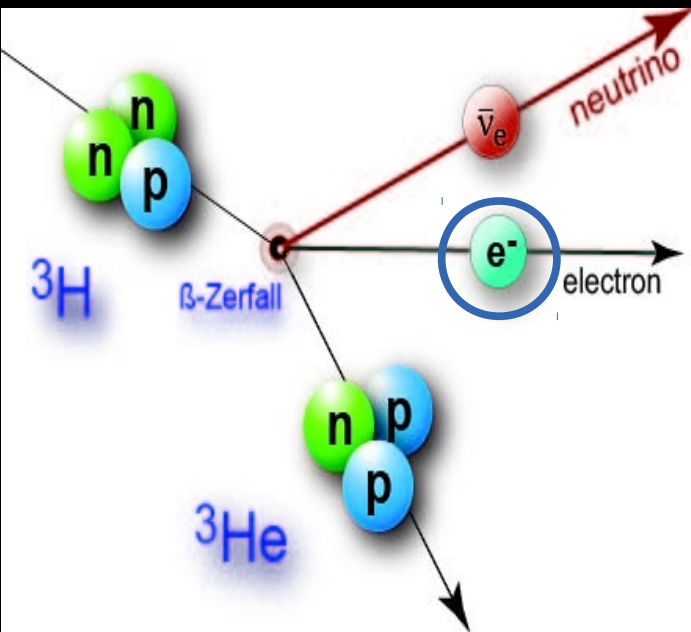
$$\frac{dR}{dE}(E) = N \frac{G_F^2}{2\pi^3 \hbar^7 c^5} \cos^2(\theta_C) |M|^2 F(E, Z') \cdot$$

$$p \cdot (E + m_e c^2) \cdot (E_0 - E) \cdot \sum_i^3 U_{ei} \sqrt{(E_0 - E)^2 - m(\nu_i)^2 c^4}$$

Def.: "electron neutrino mass"

$$m(\nu_e)^2 = \sum_i^3 |U_{ei}^2| m_i^2$$

$$\approx \sqrt{(E_0 - E)^2 - m(\nu_e)^2 c^4}$$



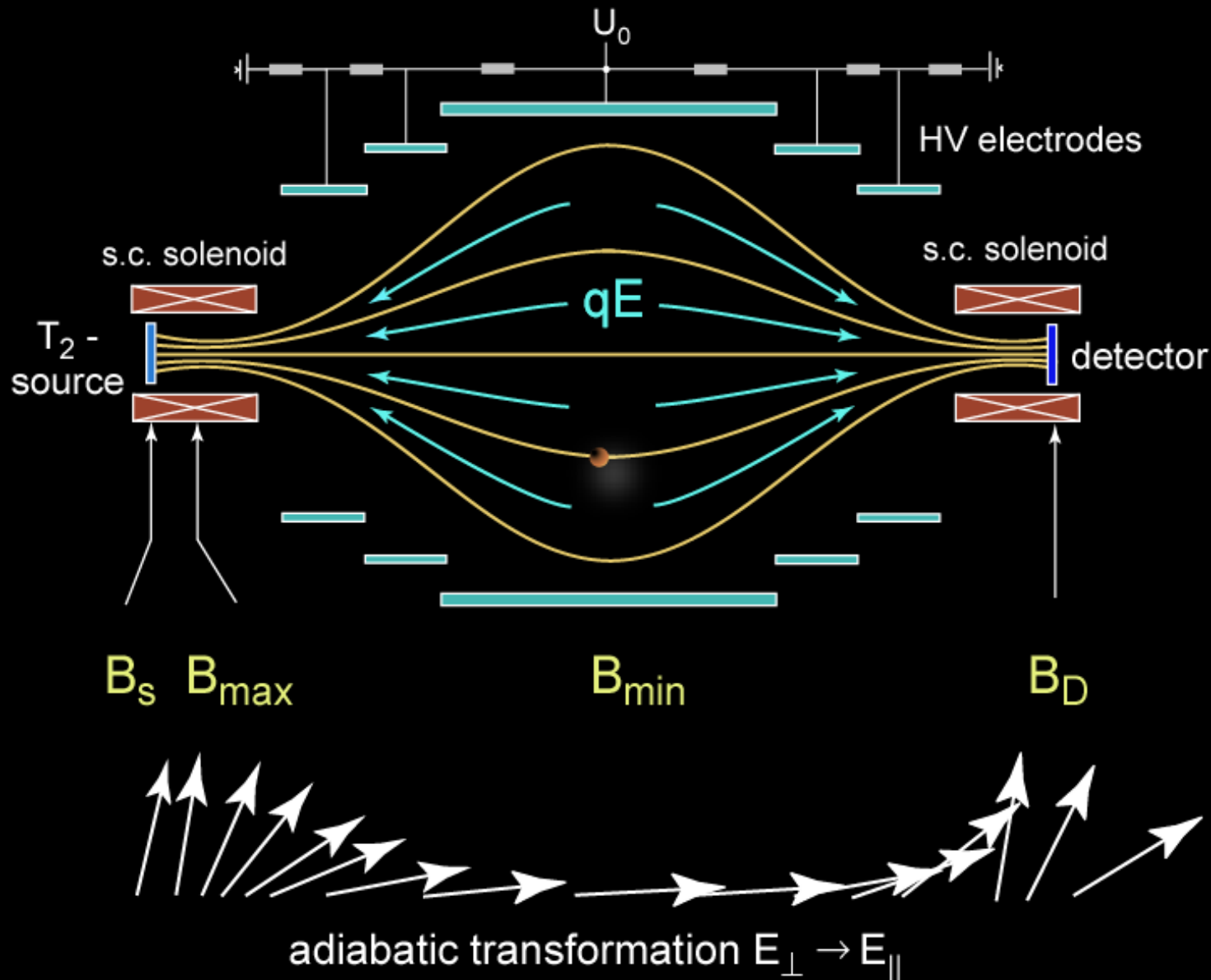
# KATRIN

Design goal:  
Measurement of  $m(\nu_e)$  with  
< 0.2 eV sensitivity (90% C.L.)

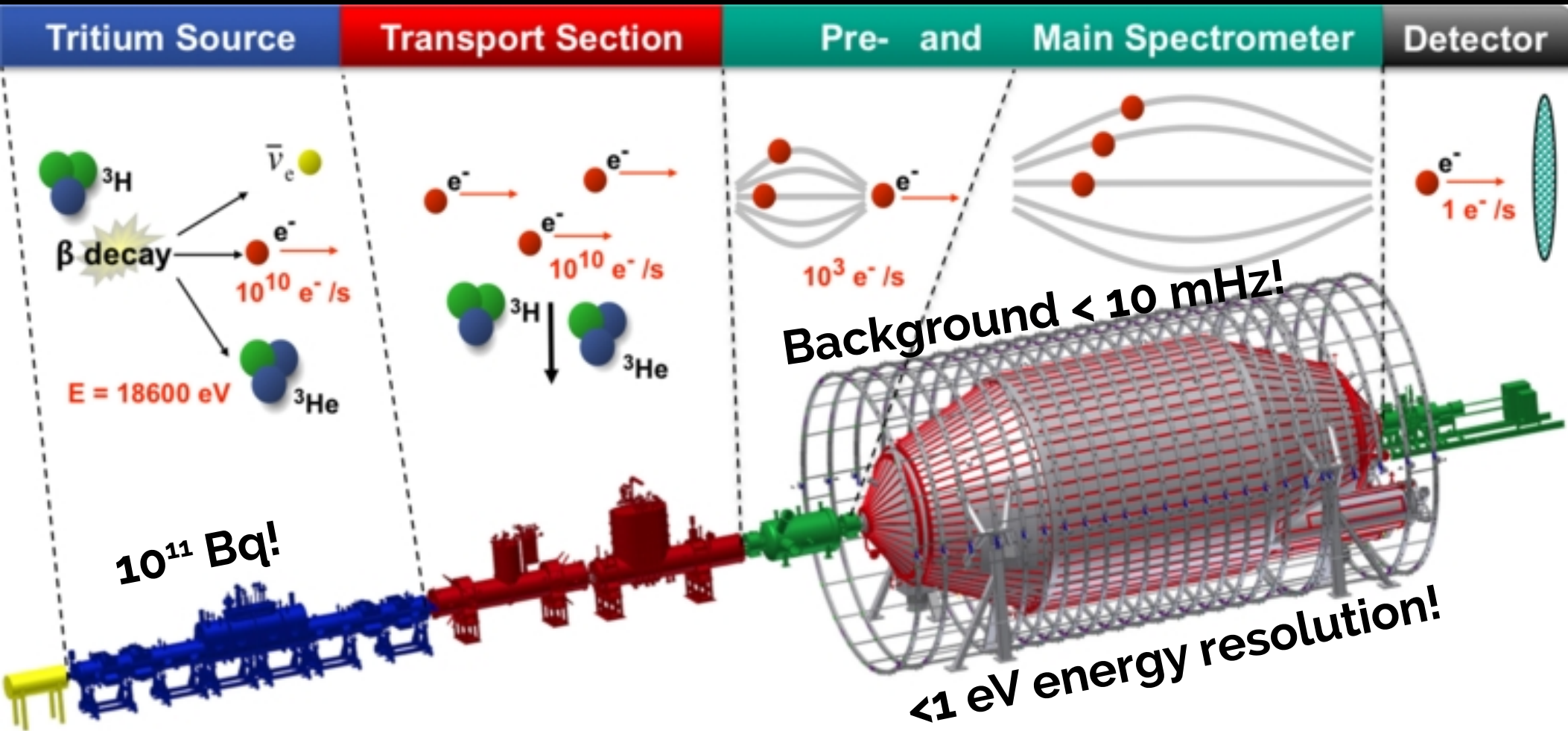




# Measure endpoint region of $\beta$ spectrum by electrostatic high pass filter (MAC-E-Filter)



# KATRIN Components



Tritium decays, releasing an electron and an anti-electron-neutrino. While the neutrino escapes undetected, the electron starts its journey to the detector.

Electrons are guided towards the spectrometer by magnetic fields. Tritium has to be pumped out to provide tritium-free spectrometers.

The electron energy is analyzed by applying an electrostatic retarding potential. Electrons are only transmitted if their kinetic energy is sufficiently high.

At the end of their journey, the electrons are counted at the detector. Their rate varies with the spectrometer potential and hence gives an integrated  $\beta$ -spectrum.



**Current status:**  
commissioning of  
main spectrometer  
and detector  
finished





# $\beta$ spectrum with 4<sup>th</sup> mass state

$$\frac{dR}{dE}(E) = N \frac{G_F^2}{2\pi^3 \hbar^7 c^5} \cos^2(\theta_C) |M|^2 F(E, Z').$$

$$p \cdot (E + m_e c^2) \cdot (E_0 - E) \cdot \sum_i^4 U_{ei} \sqrt{(E_0 - E)^2 - m(\nu_i)^2 c^4}$$

# $\beta$ spectrum with 4<sup>th</sup> mass state

$$\frac{dR}{dE}(E) = N \frac{G_F^2}{2\pi^3 \hbar^7 c^5} \cos^2(\theta_C) |M|^2 F(E, Z') \cdot$$
$$p \cdot (E + m_e c^2) \cdot (E_0 - E) \cdot \sum_i^4 U_{ei} \sqrt{(E_0 - E)^2 - m(\nu_i)^2 c^4}$$

# $\beta$ spectrum with 4<sup>th</sup> mass state

$$\frac{dR}{dE}(E) = N \frac{G_F^2}{2\pi^3 \hbar^7 c^5} \cos^2(\theta_C) |M|^2 F(E, Z') \cdot$$

$$p \cdot (E + m_e c^2) \cdot (E_0 - E) \cdot \sum_i U_{ei} \sqrt{(E_0 - E)^2 - m(\nu_i)^2 c^4}$$

$$\approx \cos \theta_4 \sqrt{(E_0 - E)^2 - m(\nu_e)^2 c^4} + \sin \theta_4 \sqrt{(E_0 - E)^2 - m(\nu_4)^2 c^4}$$

$\theta_4$ : mixing angle between  $\nu_{1/2/3}$  and  $\nu_4$

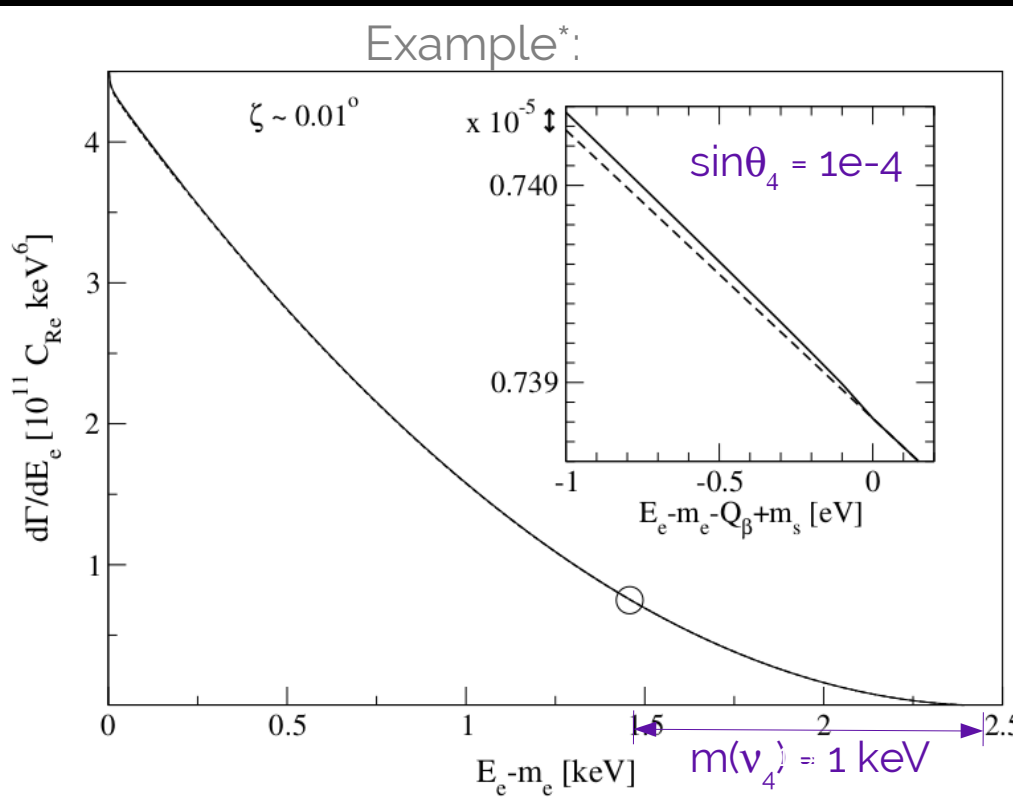
→ Linear combination of classic  $\beta$  spectrum and  $\beta$  spectrum with "heavy neutrino"



# $\beta$ spectrum with 4<sup>th</sup> mass state

$$\frac{dR}{dE}(E) = N \frac{G_F^2}{2\pi^3 \hbar^7 c^5} \cos^2(\theta_C) |M|^2 F(E, Z').$$

$$p \cdot (E + m_e c^2) \cdot (E_0 - E) \cdot \sum_i U_{ei} \sqrt{(E_0 - E)^2 - m(\nu_i)^2 c^4}$$



$$\approx \cos \theta_4 \sqrt{(E_0 - E)^2 - m(\nu_e)^2 c^4} + \sin \theta_4 \sqrt{(E_0 - E)^2 - m(\nu_4)^2 c^4}$$

$\theta_4$ : mixing angle between  $\nu_{1/2/3}$  and  $\nu_4$

→ Linear combination of classic  $\beta$  spectrum and  $\beta$  spectrum with "heavy neutrino"

\*) for  $^{187}\text{Re} \rightarrow ^{187}\text{Os}$ ,  
de Vega et al., Nucl. Phys. B866 177-195 (2013), arXiv: 1109.3452

**Can we detect keV sterile  $\nu$  with  
KATRIN?**

# Can we detect keV sterile $\nu$ with KATRIN?

- Signal very weak:  $\sin^2 \theta_4 < 10^{-7}$  (Chandra X-Ray  $\nu_s$  decay)

# Can we detect keV sterile $\nu$ with KATRIN?

- Signal very weak:  $\sin^2 \theta_4 < 10^{-7}$  (Chandra X-Ray  $\nu_s$  decay)
- Heavy systematics, e.g.:
  - Electronic final states: 43 % of molecules decay into excited states
  - Inelastic scattering: 58 % of electrons lose energy in collisions!

# Can we detect keV sterile $\nu$ with KATRIN?

- Signal very weak:  $\sin^2 \theta_4 < 10^{-7}$  (Chandra X-Ray  $\nu_s$  decay)
- Heavy systematics, e.g.:
  - Electronic final states: 43 % of molecules decay into excited states
  - Inelastic scattering: 58 % of electrons lose energy in collisions!
- High pass filter not suited to measure at  $\sim$  keV below  $\beta$  endpoint



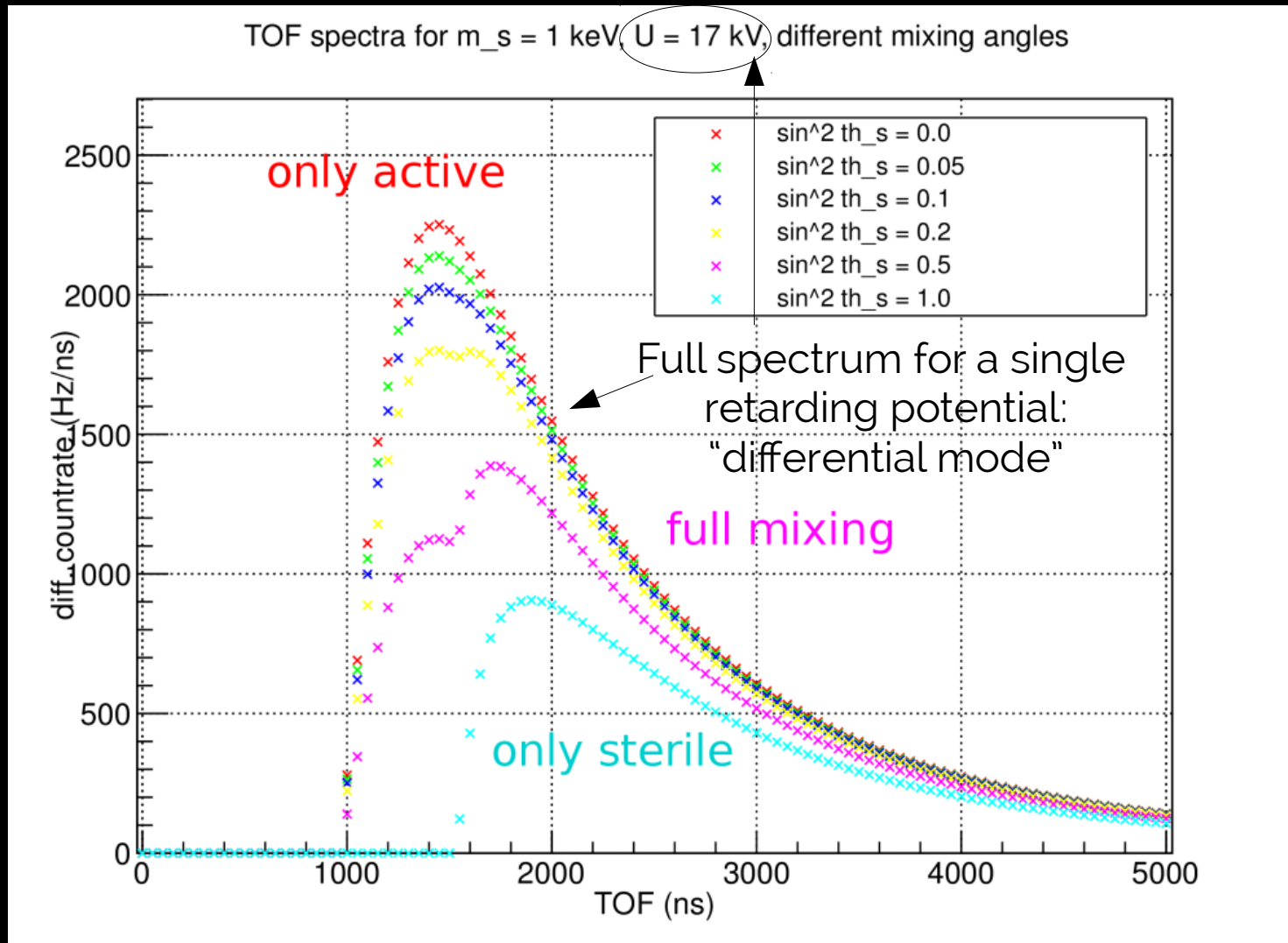
# Can we detect keV sterile $\nu$ with KATRIN?

- Signal very weak:  $\sin^2 \theta_4 < 10^{-7}$  (Chandra X-Ray  $\nu_s$  decay)
- Heavy systematics, e.g.:
  - Electronic final states: 43 % of molecules decay into excited states
  - Inelastic scattering: 58 % of electrons lose energy in collisions!
- High pass filter not suited to measure at  $\sim$  keV below  $\beta$  endpoint

=> **New strategy necessary!**

# Idea: Time-Of-Flight (TOF) mode

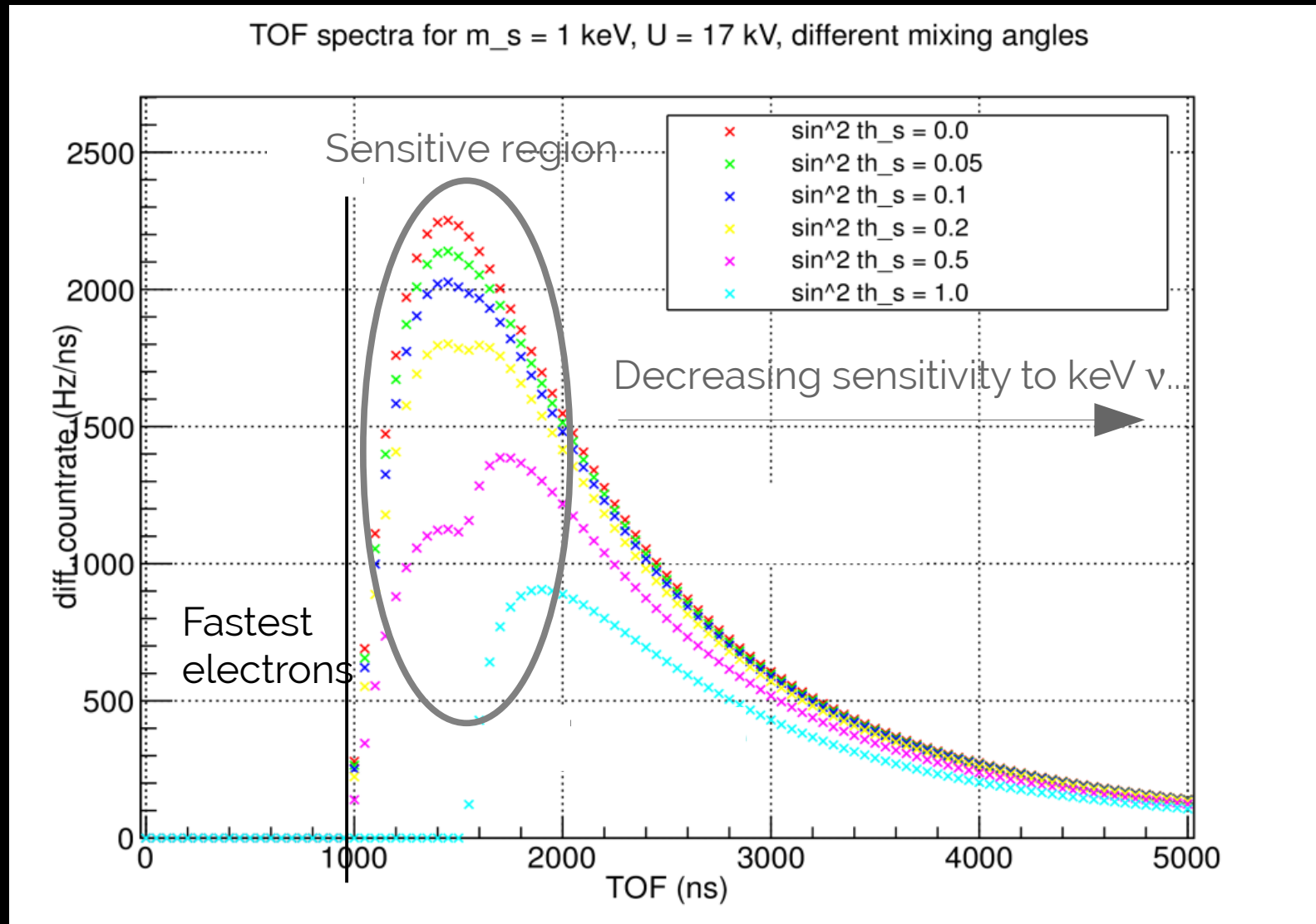
TOF function of energy  $\Leftrightarrow$  TOF spectrum function of  $\beta$  spectrum



For active  $\nu$ : potential mass sensitivity improvement (stat) up to factor 2!  
(*NSt, Christian Weinheimer et al., 2013, arXiv:1308.0532, submitted to NJP*)

# 'Band pass' like behavior of TOF

Can distinguish  $v_4$  sensitive parts from higher energy parts!



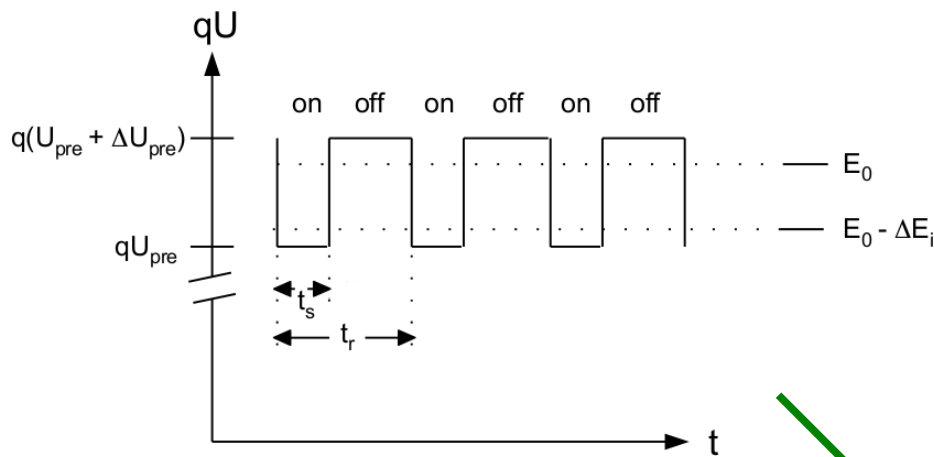
# TOF measurement: gated filter

- 1)** Pulse pre-spectrometer potential with narrow duty cycle  
→  $e^-$  pass during 'on' cycles

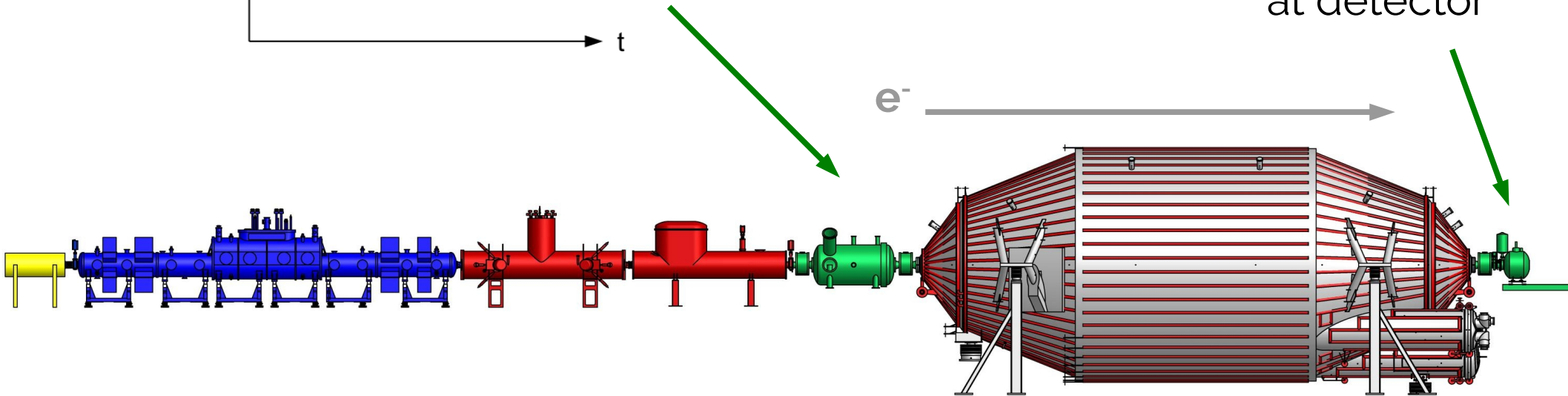
Disadvantage:

Loss of count rate.

But statistics no problem  
(S. Mertens et al., t.b.p.)



- 2)** Measure time-of-arrival at detector

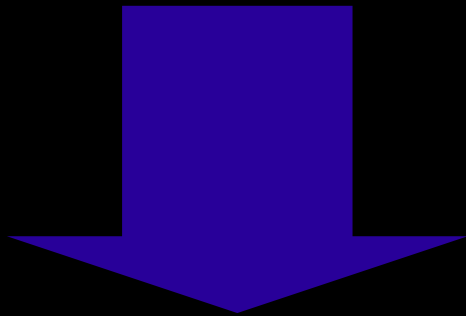


# MC simulations needed!

Systematics...

Parameter optimization...

Statistics...



Sensitivity of KATRIN-TOF to keV sterile  $\nu$

# Conclusion

- Sterile keV neutrino is DM candidate
- KATRIN may be sensitive to indirect detection
- TOF mode may improve sensitivity
- Sensitivity to be simulated



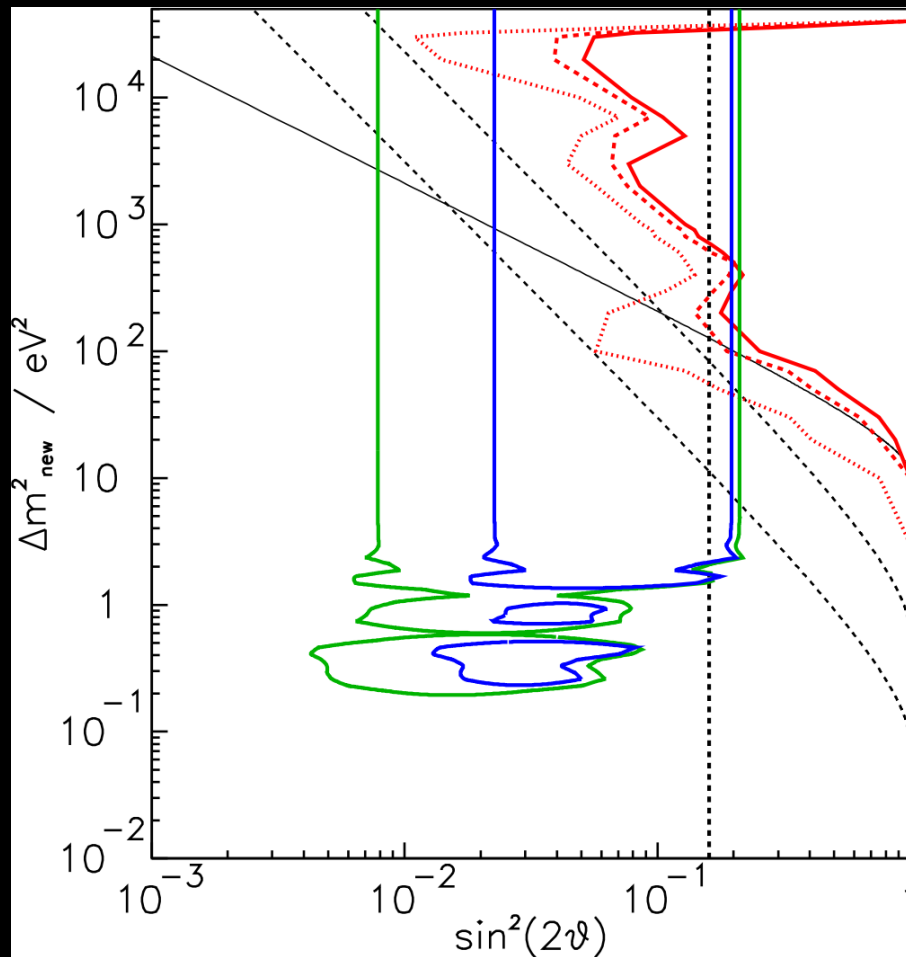
Contract no.: 05A11PM2



# Backup Slides

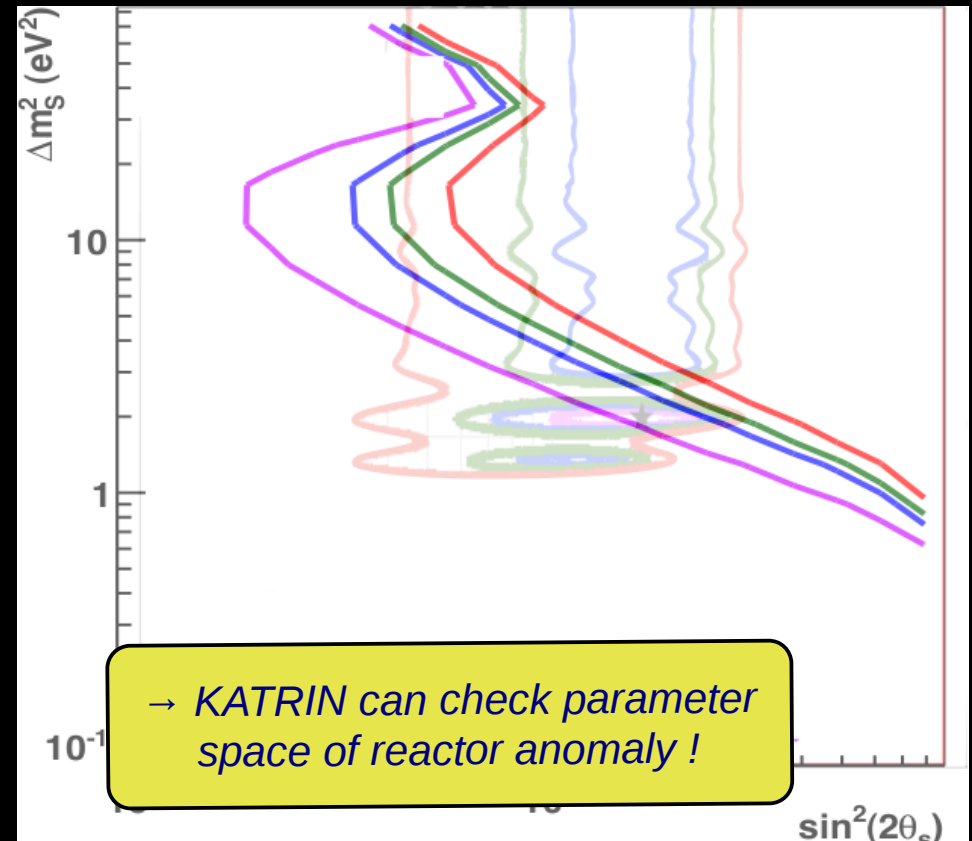
# Sensitivity to light sterile $\nu$

## MAINZ



Ch. Kraus, A. Singer, K. Valerius, C. Weinheimer, arXiv.1210.4194,  
Eur. Phys. J. C73 (2013)

## KATRIN



J. A. Formaggio, J. Barret, PLB 706 (2011) 68  
A. Seiersen Riis, S. Hannestad, JCAP02 (2011) 011  
A. Esmaili, O.L.G. Peres, arXiv:1203.2632

# Classic integrated vs. differential spectrum

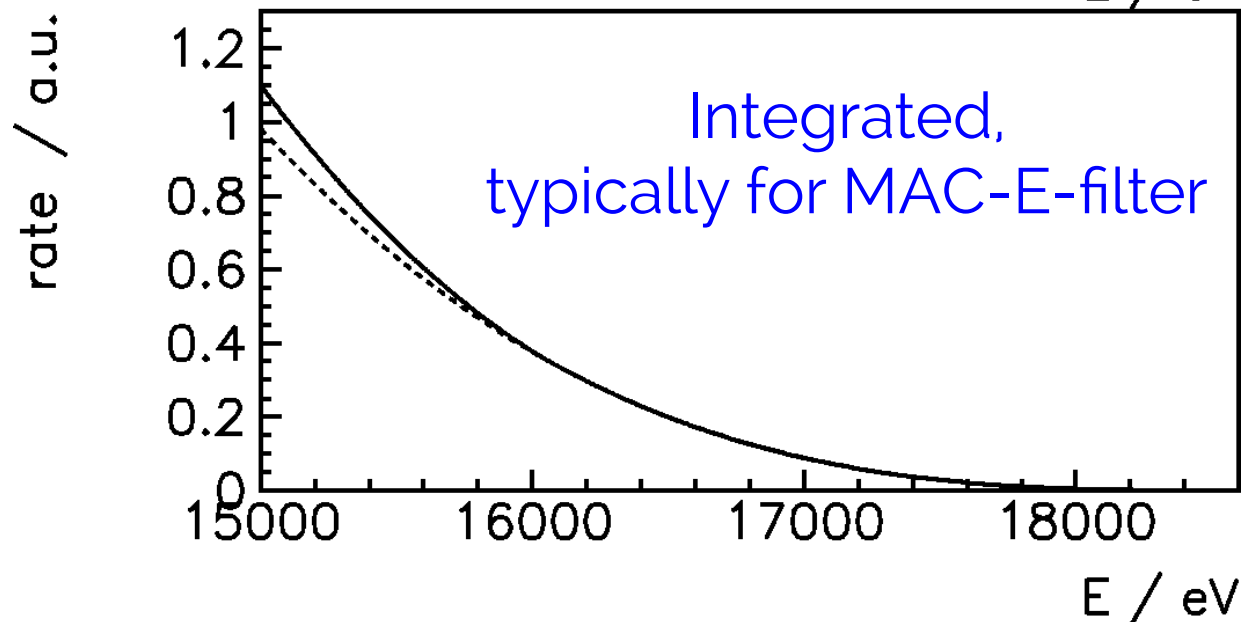
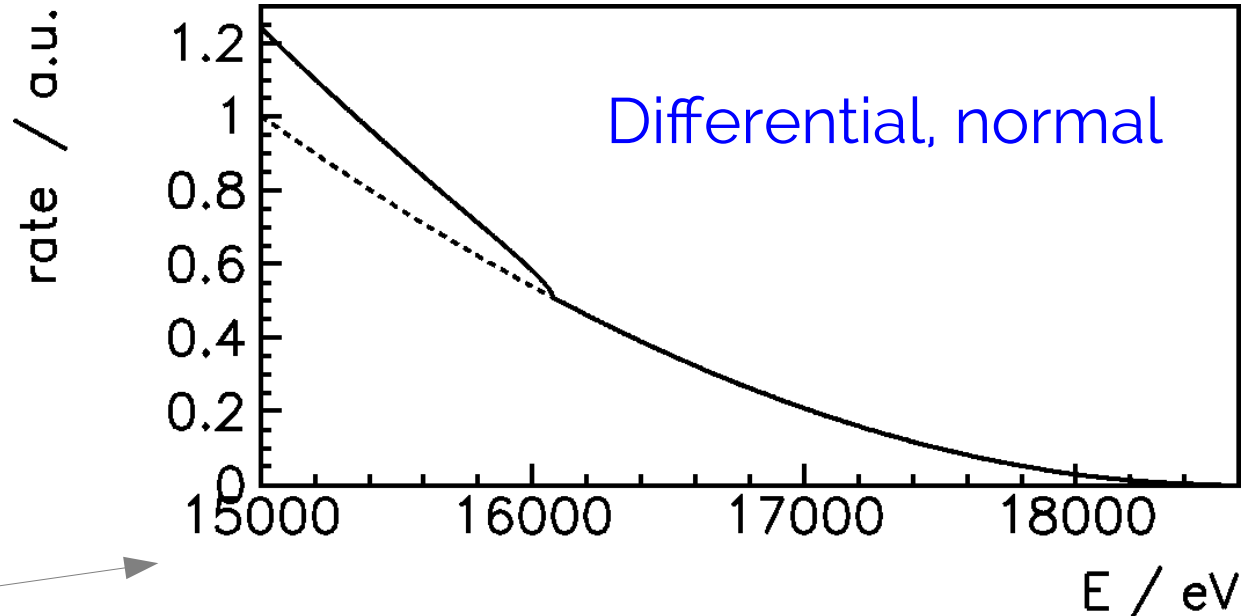
→ Obviously much better signal-to-background-ratio for differential  $\beta$ -spectrum w.r.t. integrated  $\beta$ -spectrum

Example:

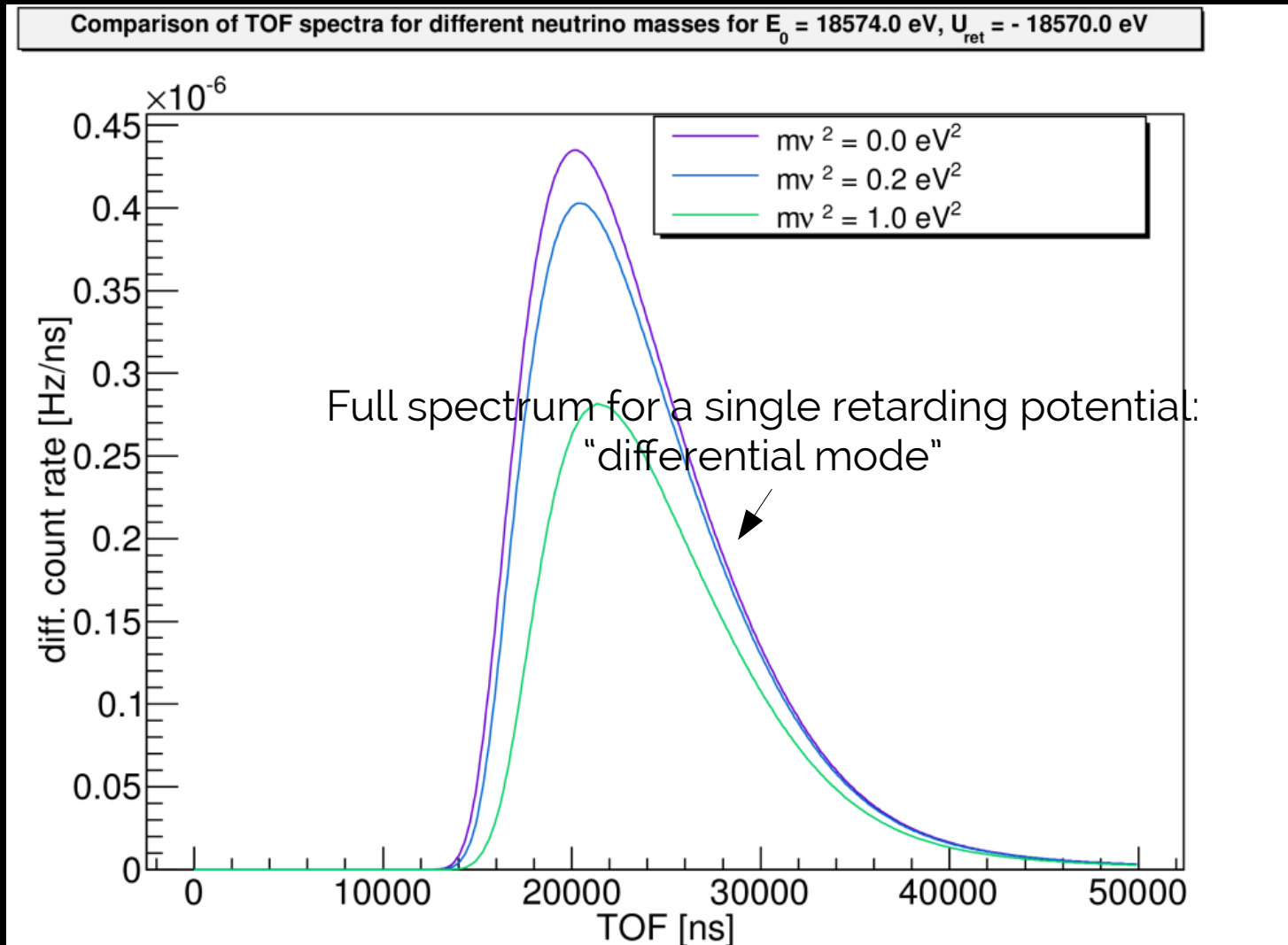
$$m(v_s) = 2.5 \text{ keV}$$

$$\sin^2(\theta) = 0.25$$

(unrealistic!)



# Neutrino mass dependence of TOF mode



Potential mass sensitivity improvement (stat) up to factor 2!  
(NSt, Christian Weinheimer et al., 2013, arXiv:1308.0532, submitted to NJP)